

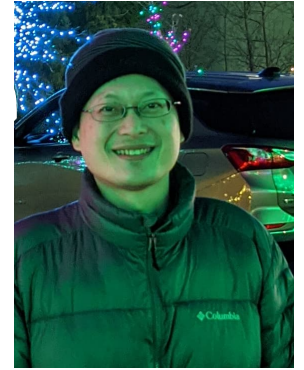


UNIVERSITY OF
MARYLAND

Jefferson Lab Accelerator Seminar

Thursday, May 27, 2021

11:00 AM – 12:00 PM (EDT)



Loss Mechanisms on Superconducting Quantum Devices and Microwave Microscopy for Probing Superconducting Devices Metrology

Speaker: Tamin Tai

Acknowledgement:

Dr. Steven Anlage Group: Dr. Behnood Ghamsari, Jingnan Cai, Dr. Bakhron Oripov

Dr. Ben Palmer Group: Dr. S. Premaratne, Dr. Jen-Hao Yeh, Rui R. Zhang, Yizhou Huang

Dr. Frederick Wellstood Group: Dr. Sudeep Dutta, Cody Ballard



U.S. DEPARTMENT OF
ENERGY

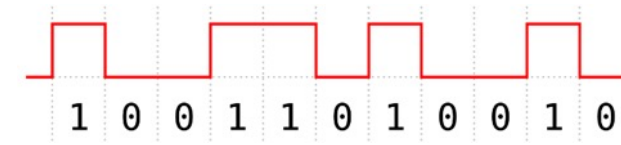


DEPARTMENT OF
PHYSICS

This work is primarily supported by DOE-BES (grant # DESC0018788)
and University of Maryland, Quantum Materials Center

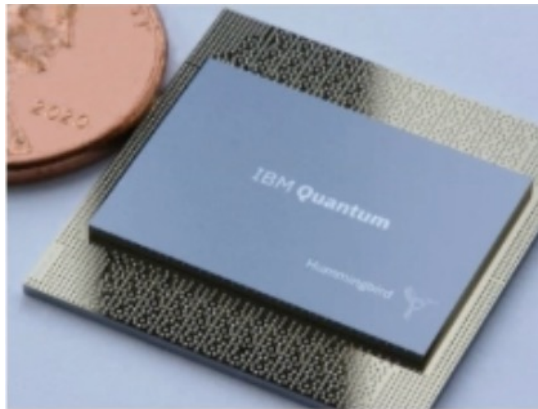
Classical vs. Quantum computing

- Classical computing utilizes two voltage levels for computations (bits)
- Quantum superpositions can be used to speed up certain computations

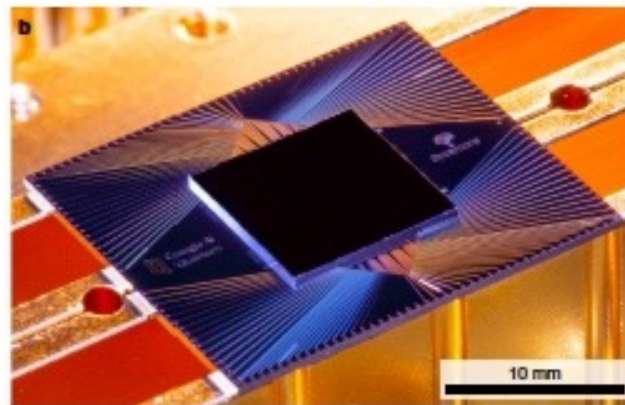


$$\Psi = \alpha|0\rangle + \beta|1\rangle$$

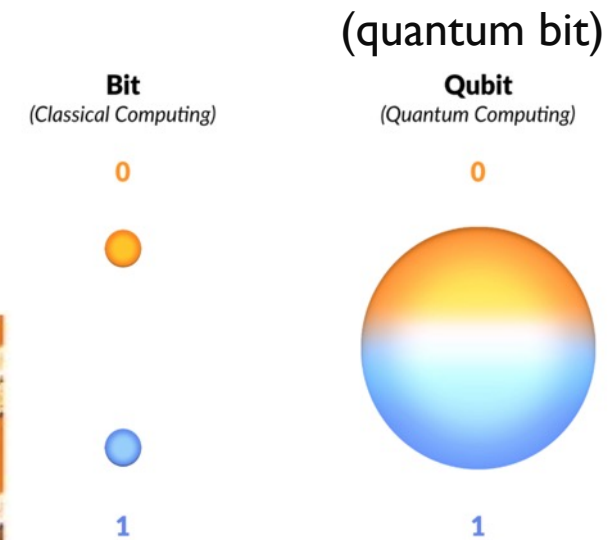
Superconducting quantum systems



IBM chip



Google



Different Types of Superconducting Qubit

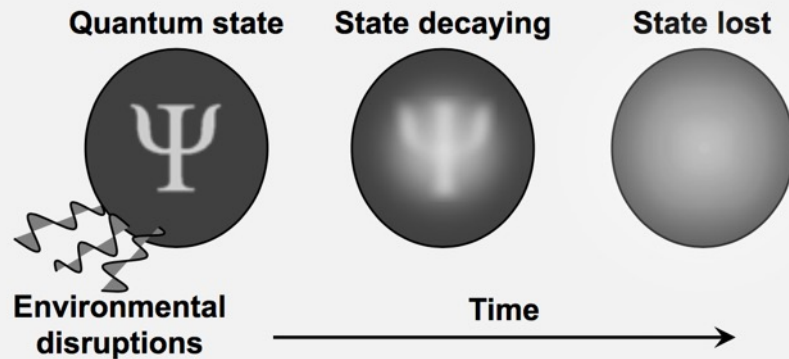
Remarkable improvement in T_1 and T_2

--- Materials

--- Design

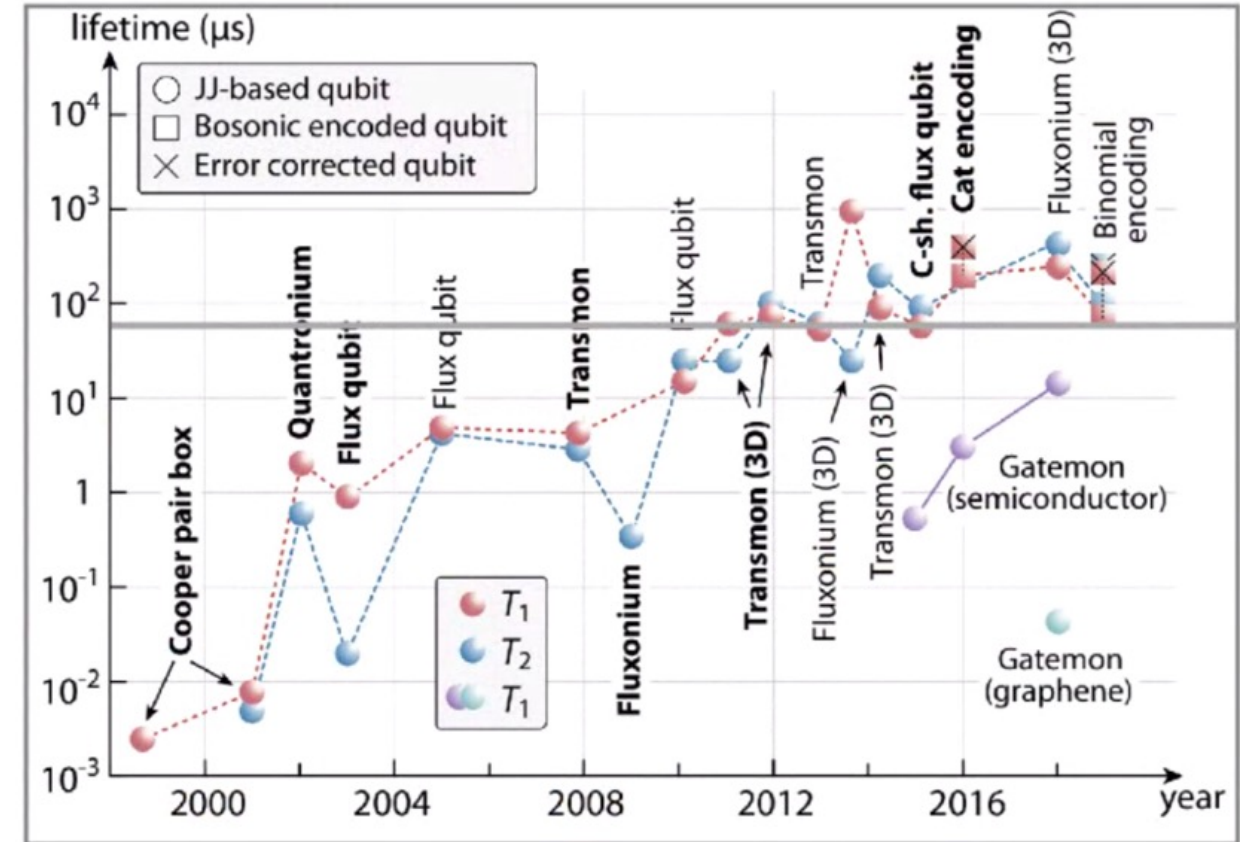
--- Fabrication

Coherence time t_{coh} : The qubit's lifetime



$$\frac{1}{T_2} = \frac{1}{2T_1} + \frac{1}{T_\phi}$$

Moore's law for T_2



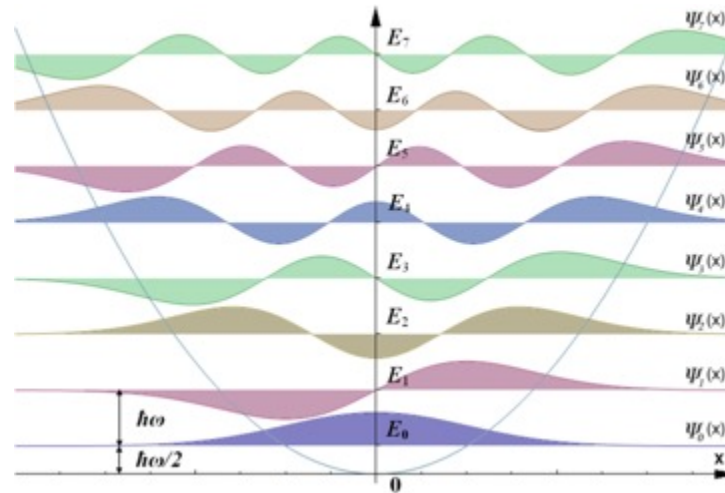
Kjaergaard *et al*, arXiv:1905.13641v1

M. Kjaergaard, WDO, et al., arXiv:1905.13641

P. Krantz, WDO, et al., Appl. Phys. Rev. 6, 021318 (2019); arXiv:1905.13641

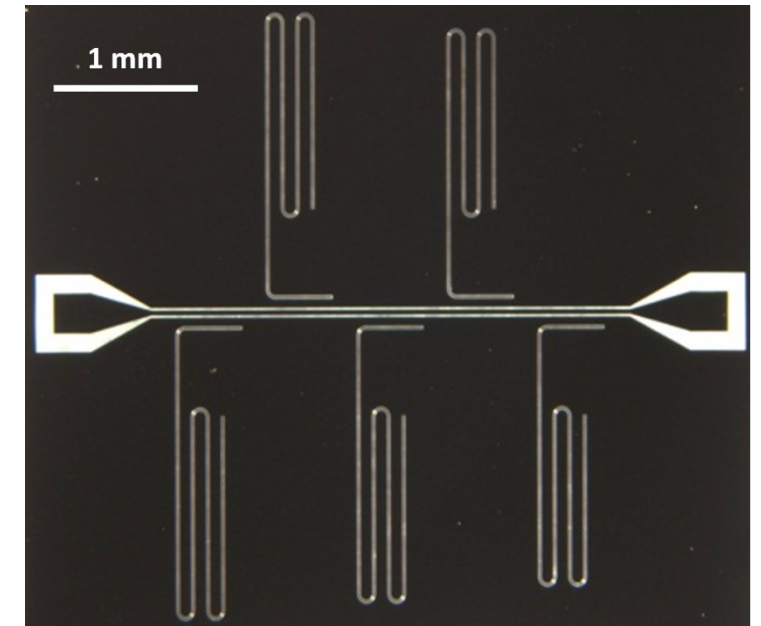
WDO & Welanders, MRS Bulletin (2013)

Harmonic Resonators



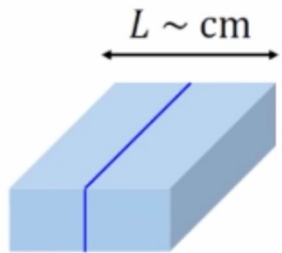
$$E_n = \left(n + \frac{1}{2}\right)\hbar\omega$$

Planar types of Cavities



G. Calusine *et al.* *Appl. Phys. Lett.*,
vol. 112, no. 6, pp. 1–7, 2018

3D types of Cavities



Ep In/Al rectangular TE₁₁₀



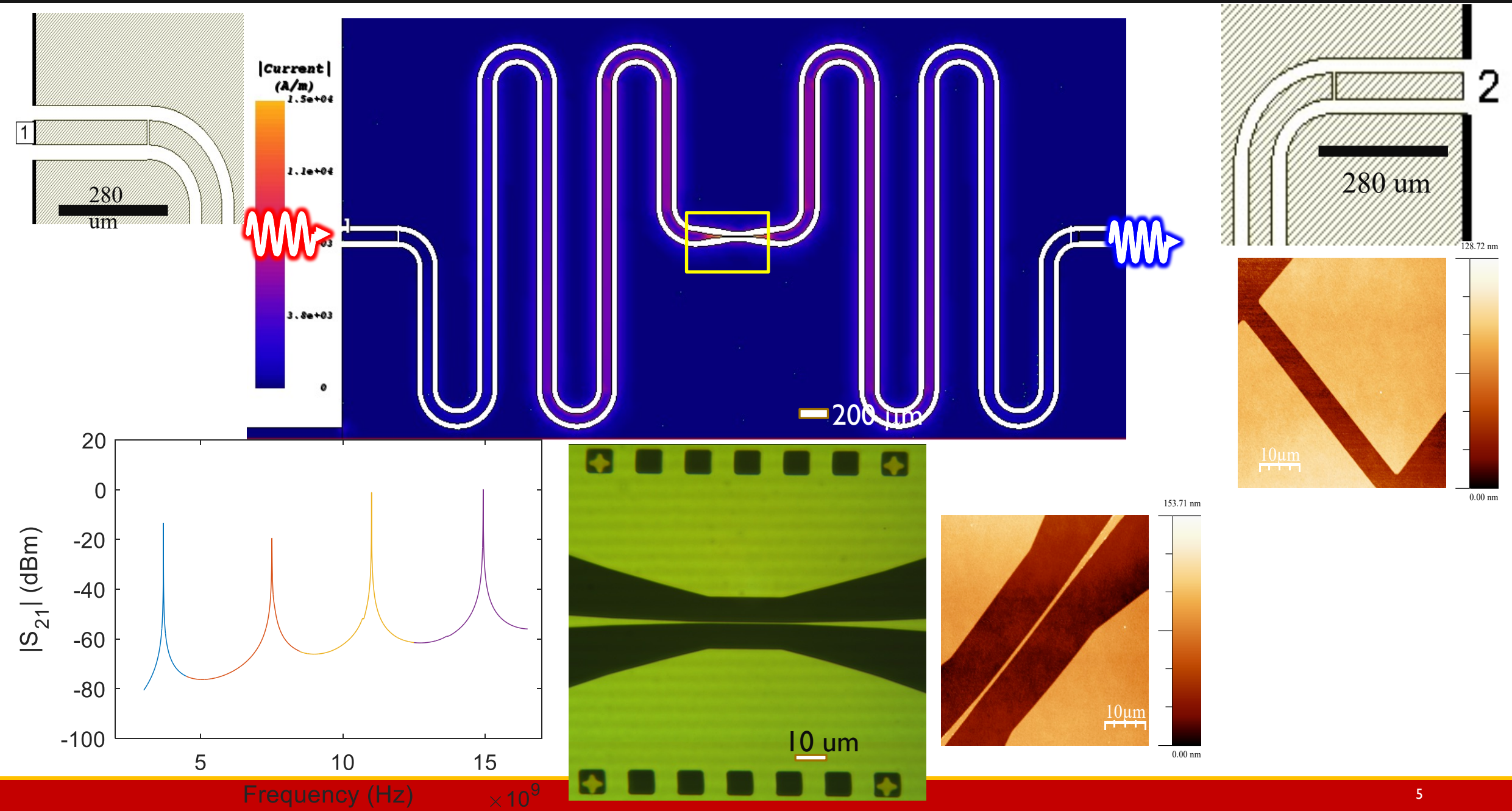
Al stub $\lambda/4$

Nb stub $\lambda/4$ [2]

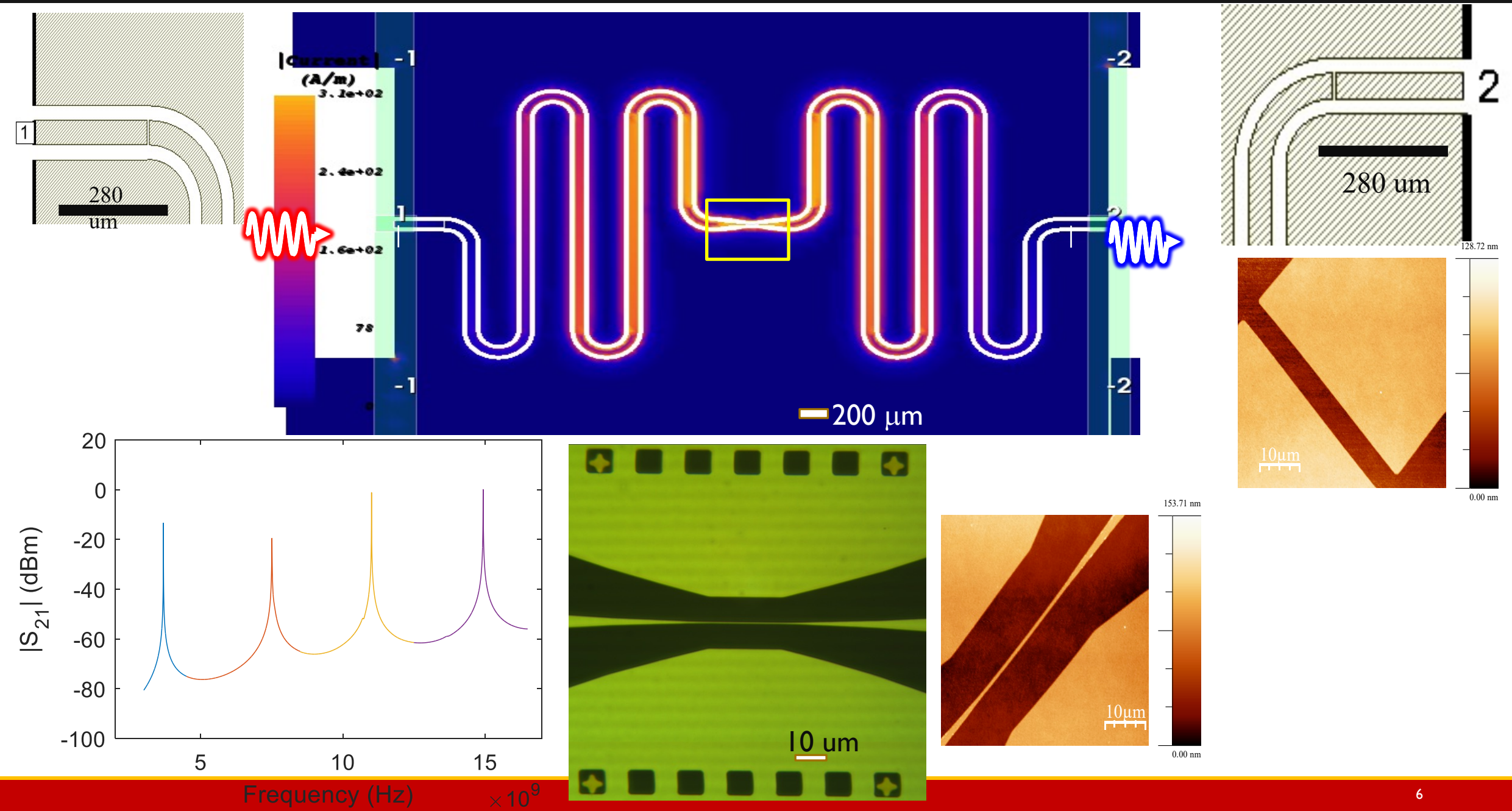


Nb TESLA TM₀₁₀ [1]

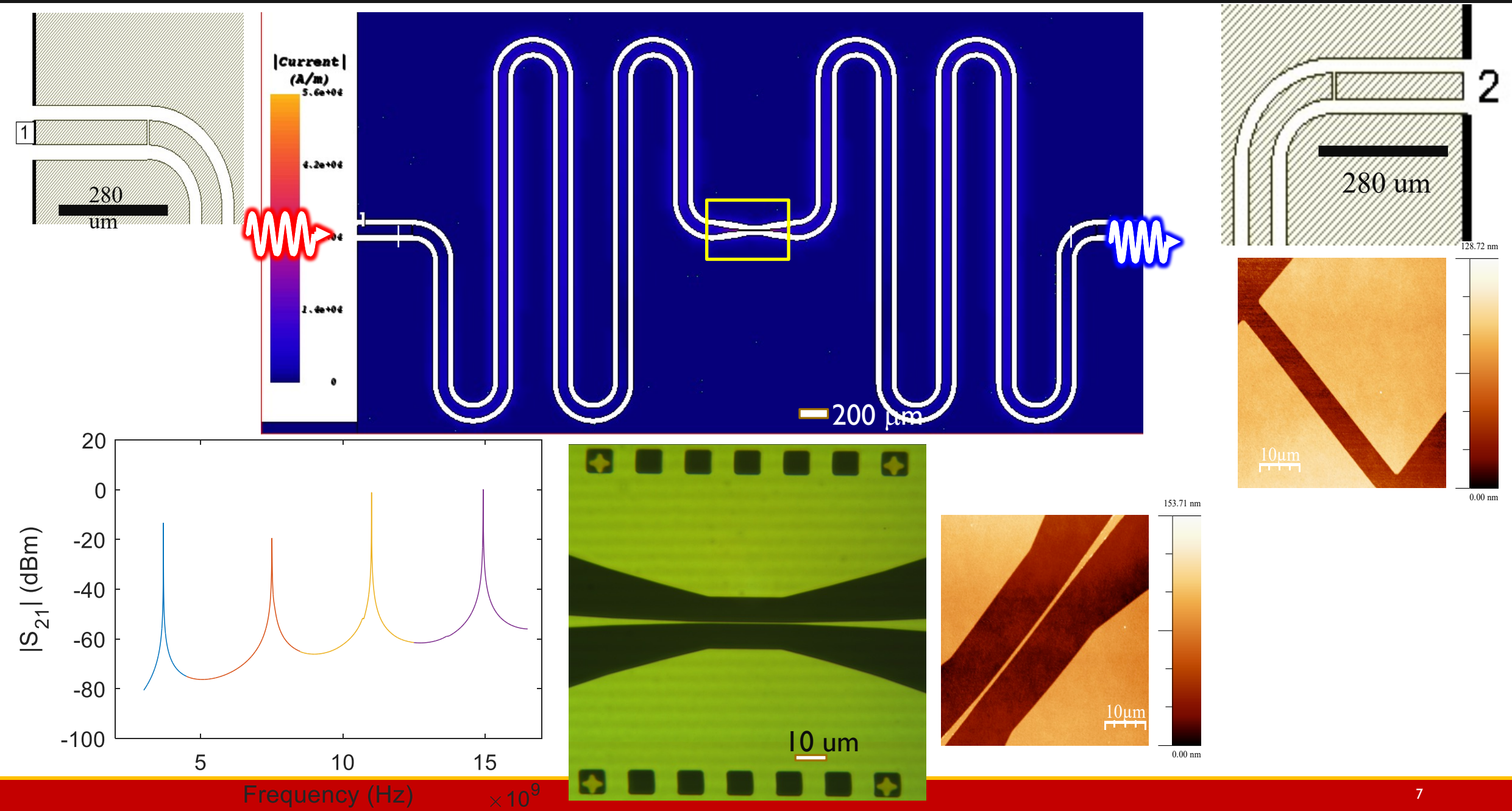
$\lambda/2$ Resonator Design and Fabrication



$\lambda/2$ Resonator Design and Fabrication



$\lambda/2$ Resonator Design and Fabrication

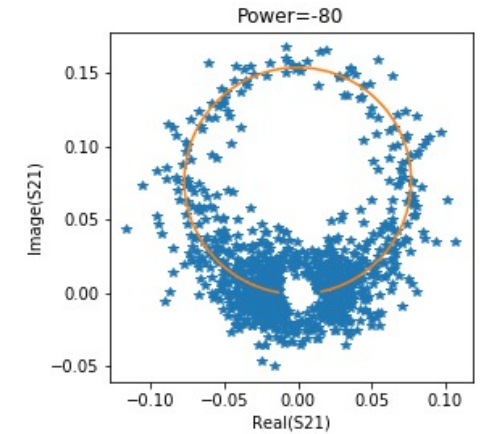
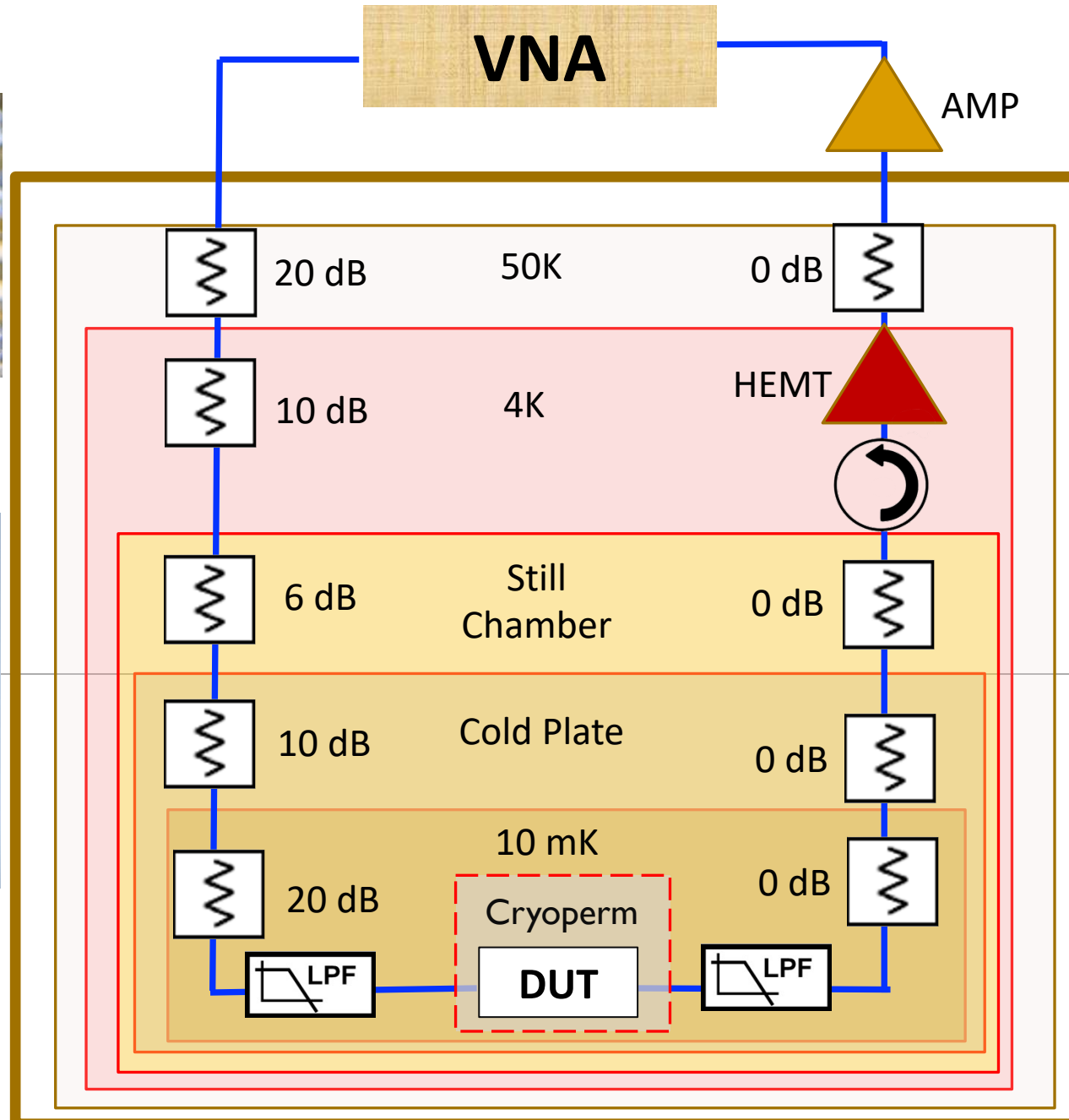
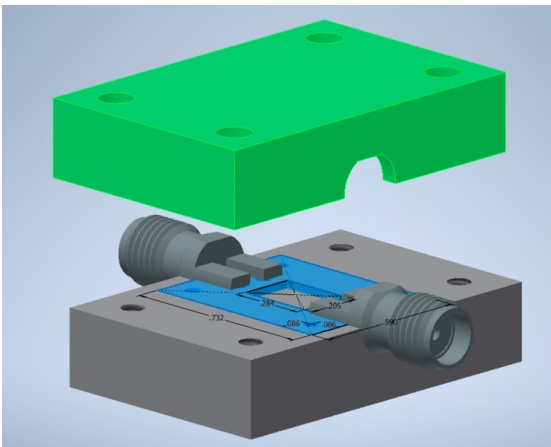


Packaging and System Setup

Indium GND



Packaging



$$S_{21}(f) = |S_{21,in}| |S_{21,out}| * \left(\frac{Q_L/Q_c}{1 + 2iQ_L \left(\frac{f-f_r}{f_r} \right)} + c_0 \right)$$

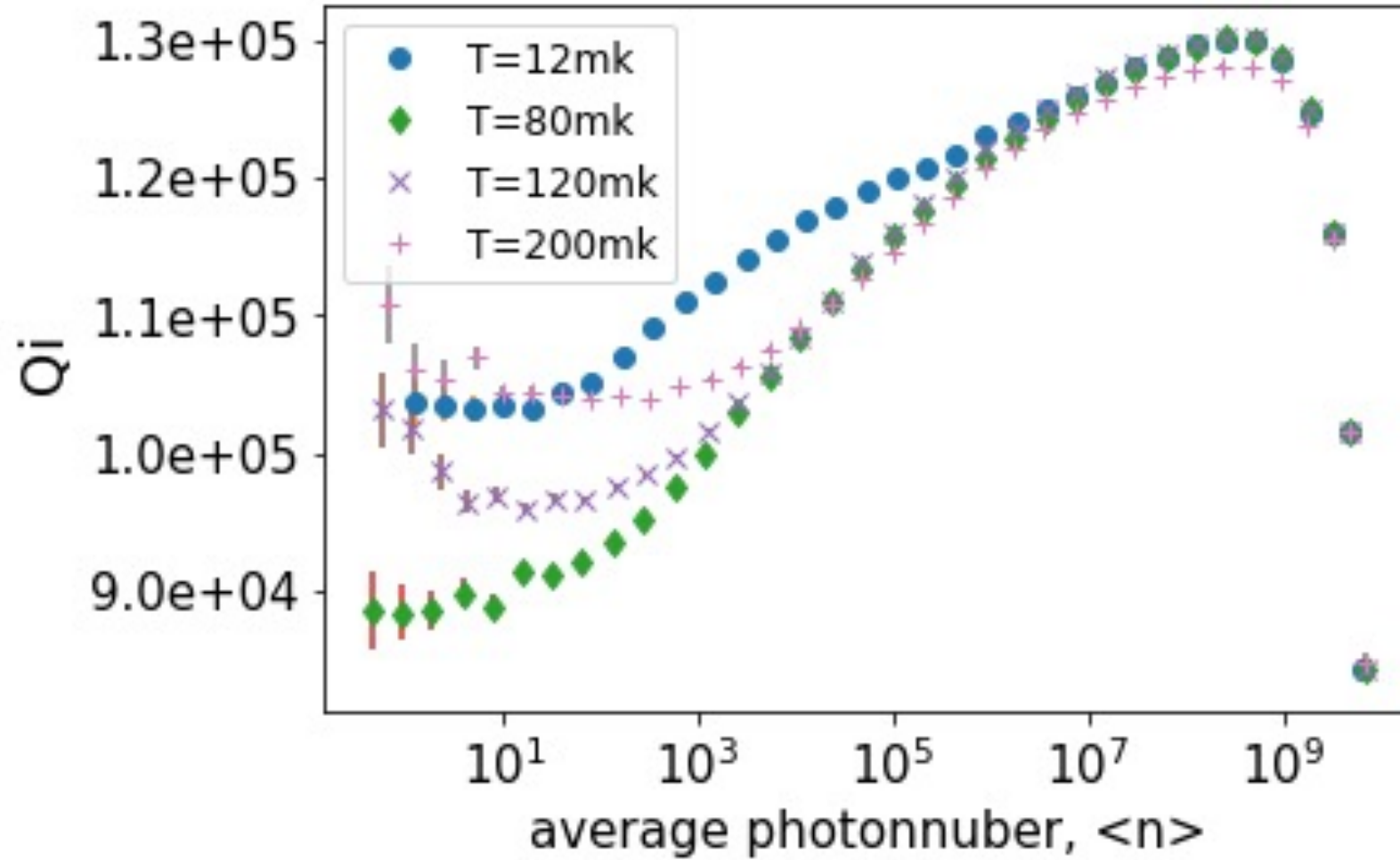
$$\Rightarrow Q_L \text{ and } Q_c$$

$$\frac{1}{Q_L} = \frac{1}{Q_i} + \frac{1}{Q_c}$$

$$Q_c \sim 7e5$$

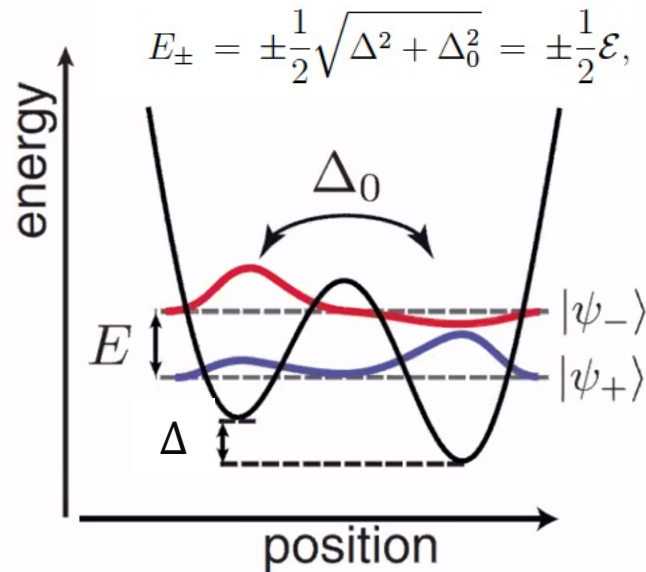
$$Q_L \sim Q_i$$

Measurement Results

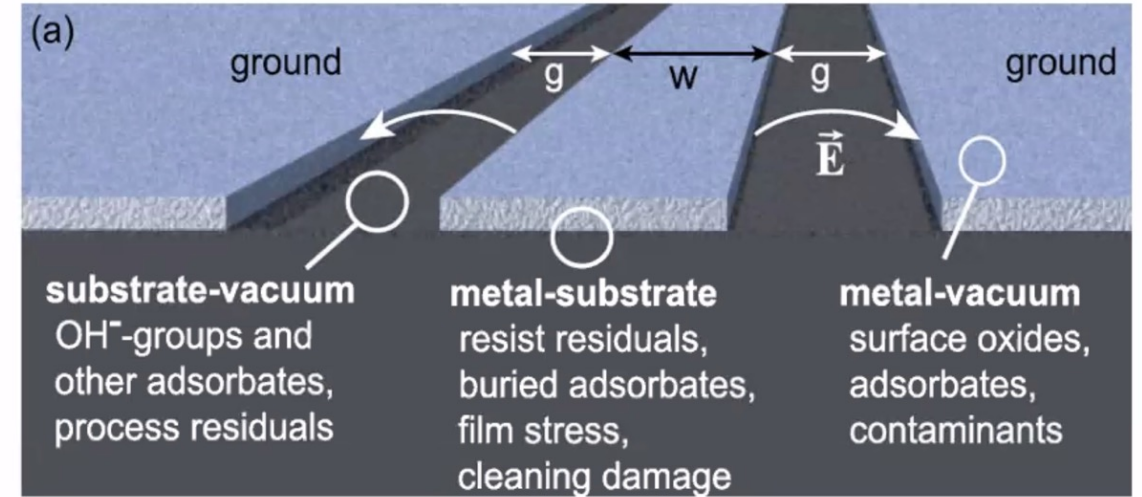


$$\langle n \rangle \equiv \frac{4Q_L^2 P_{in}}{Q_c \hbar \omega_0^2}$$

Two Level System Losses



Schematic of a Two Level System²



Muller Rep. Prog. Phys. (2019)

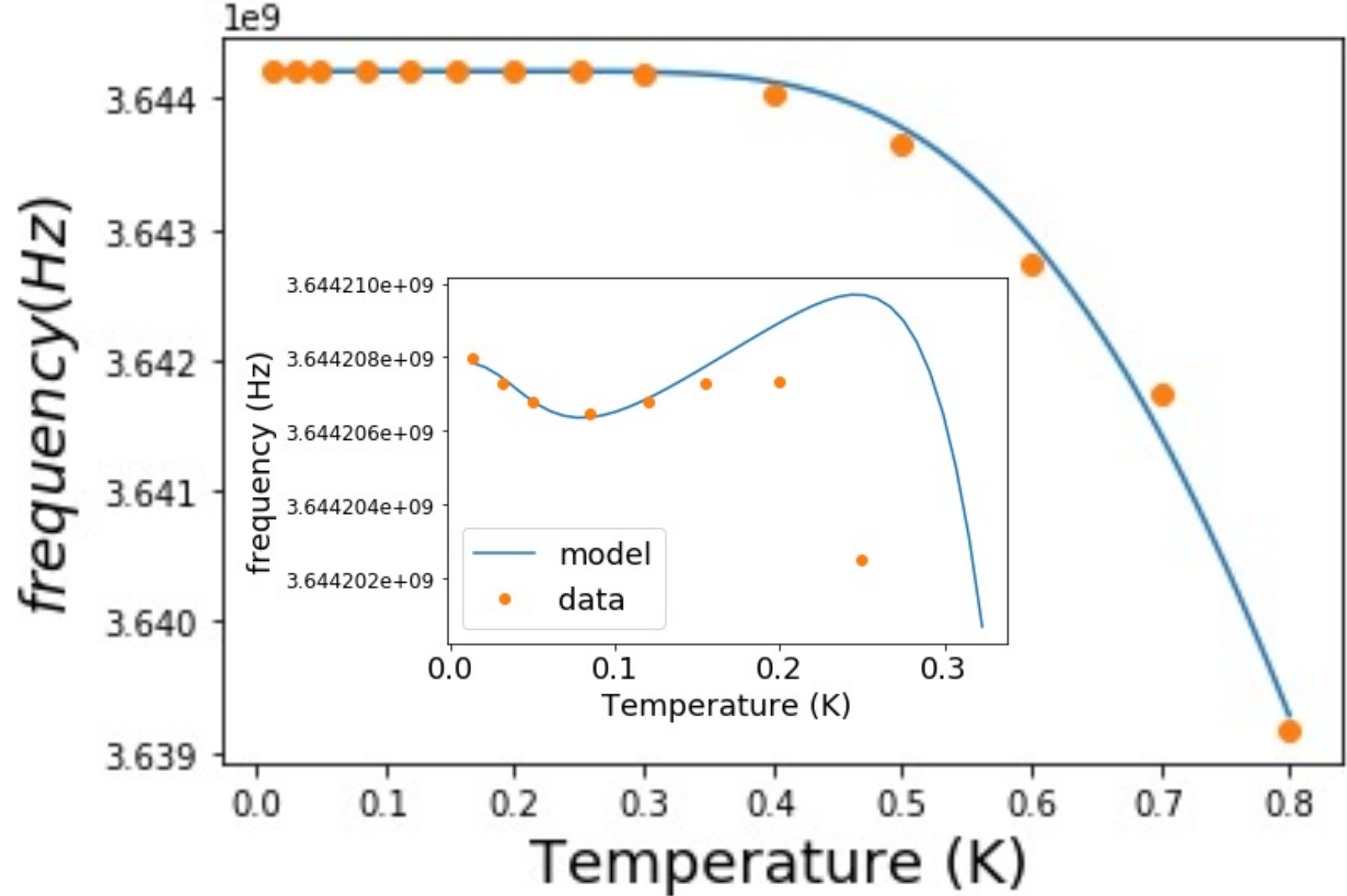
• Sources of TLS

1. Defects present in dielectrics at the metal/substrate/vacuum interface
2. Primarily associated with oxide layer/OH- groups/chemical residues

$$H_{\text{TLS}} = \begin{pmatrix} -\Delta & \Delta_0 \\ \Delta_0 & \Delta \end{pmatrix}$$

Contribution of TLS and Thermal Quasi Particles

$$\frac{f-f_n}{f_n} = F * \frac{\tan\delta}{\pi} \left\{ \operatorname{Re} \left[\Psi \left(0.5 - \frac{hf}{2\pi j k_b T} \right) \right] - \ln \left(\frac{hf}{k_b T} \right) \right\} - \frac{\alpha}{2} \sqrt{\frac{\pi \Delta_{s0}}{2k_b T}} \exp \left(-\frac{\Delta_{s0}}{k_b T} \right)$$

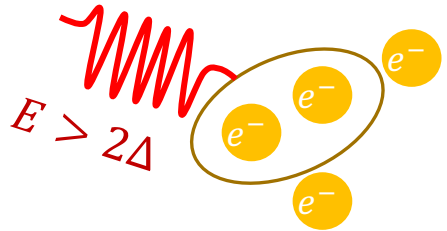


$$\alpha = \frac{L_k}{L_k + L_{geo}}$$

Fitting	Value
Δ_{s0} (μeV)	174.03 μeV
α	0.0185
F	0.32
$\tan\delta$	2e-5 (sapphire)
f_n (GHz)	3.644

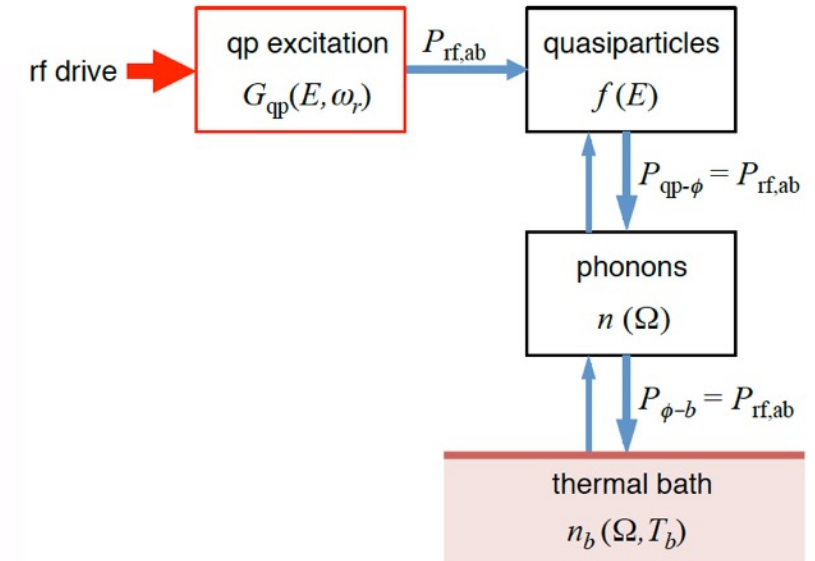
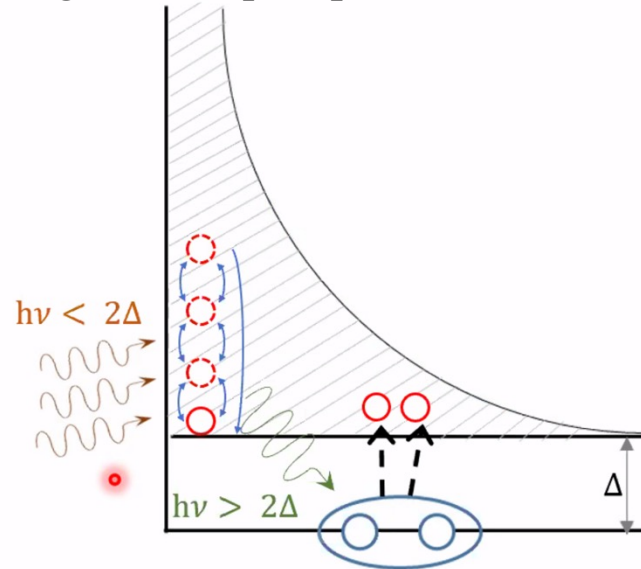
Contribution of Nonequilibrium Quasiparticles and Others

Δ - superconducting energy gap



$E > 2\Delta$ breaks a Cooper pair forming 2 quasiparticles giving dissipation

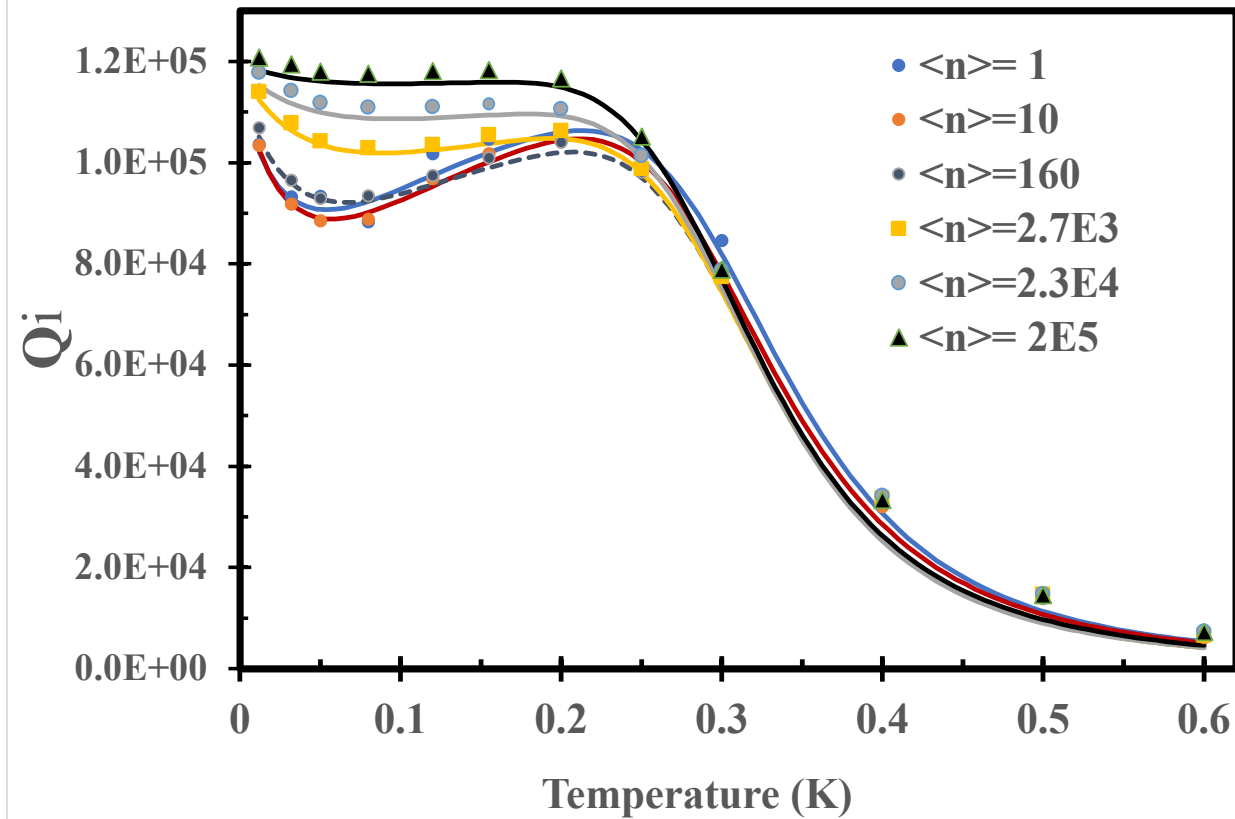
Multi photon $h\nu < 2\Delta$ also generate quasiparticles



$$\frac{1}{Q_i} = \frac{1}{Q_{TLS}(T, E)} + \frac{1}{Q_{QP}(T, E)} + \frac{1}{Q_{vortices}(B)} + \frac{1}{Q_{rad}(w)}$$

A diagram of a vortex core, represented by a dashed circle with a central electron (e^-) and the word "vertex" above it.

Temperature Dependent Q_i



$$\delta_{TLS} = \frac{\delta_{TLS}^0 \tanh\left(\frac{\varepsilon}{2k_B T}\right)}{\sqrt{1 + \frac{A}{T} \tanh\left(\frac{\varepsilon}{2k_B T}\right)}}$$

$$T_1 \sim \tanh\left(\frac{\varepsilon}{2k_B T}\right)$$

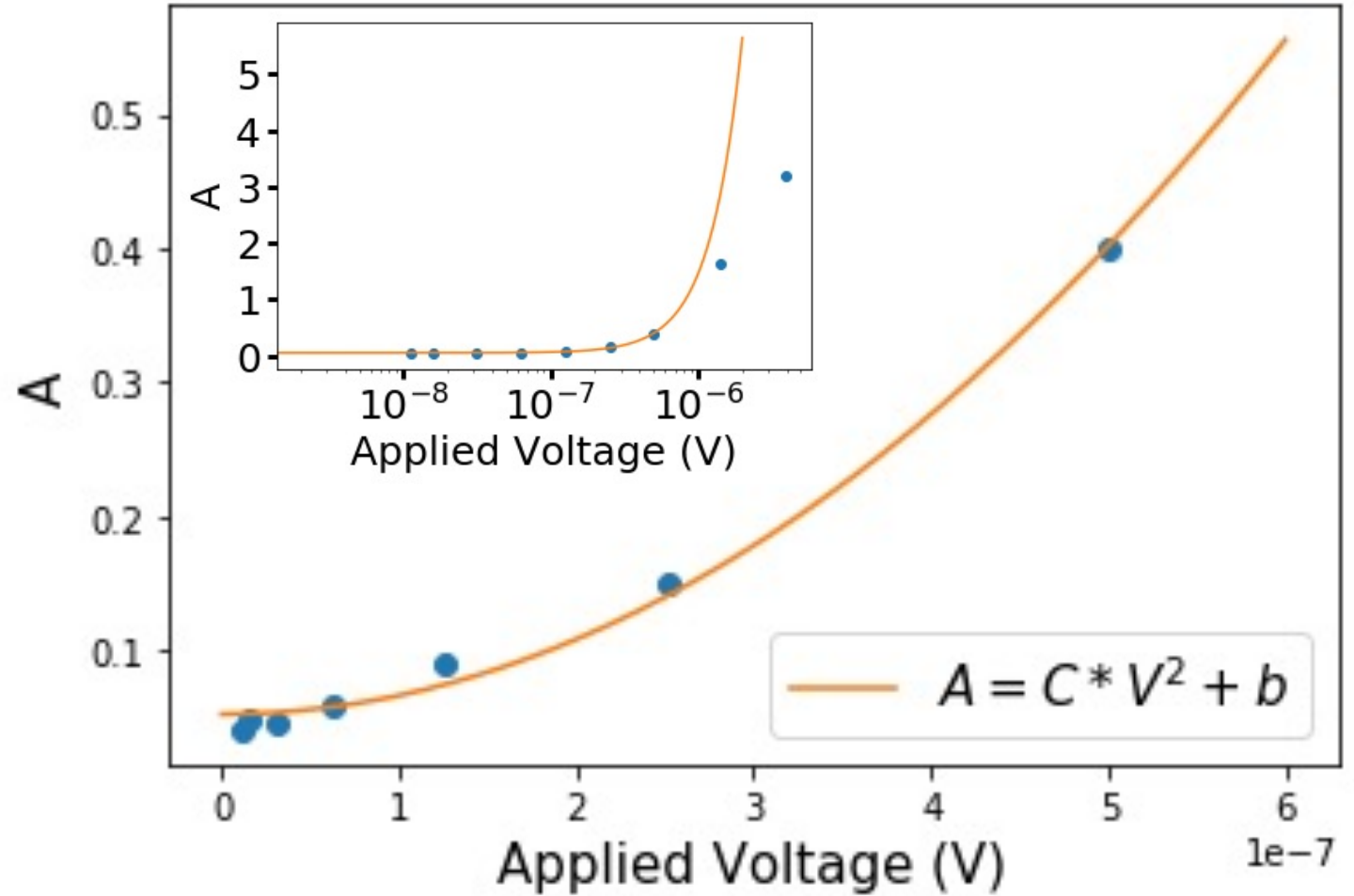
$$T_2 \sim 1/T$$

$$\frac{1}{Q_i} = \frac{1}{Q_{TLS}(T, E)} + \frac{1}{Q_{QP}(T, E)} + \frac{1}{Q_{other}}$$

Field Dependence of Fitting Parameter

$$\delta_{TLS} = \frac{\delta_{TLS}^0 \tanh\left(\frac{\varepsilon}{2k_B T}\right)}{\sqrt{1 + \frac{A}{T} \tanh\left(\frac{\varepsilon}{2k_B T}\right)}}$$

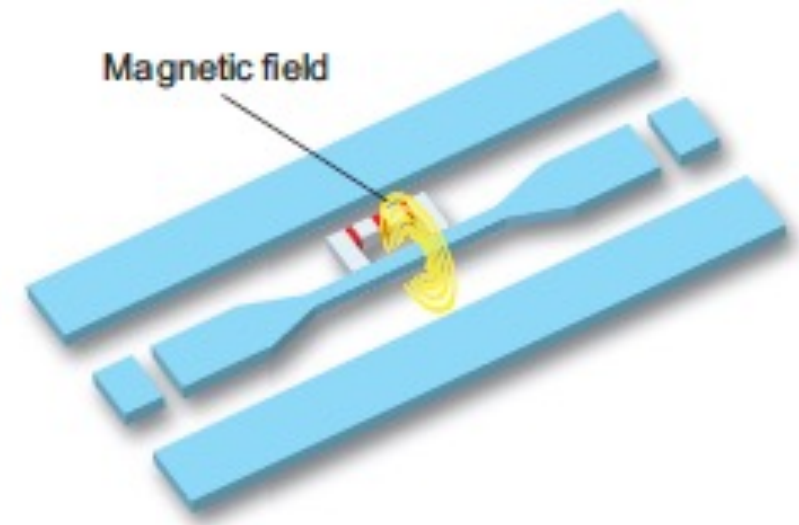
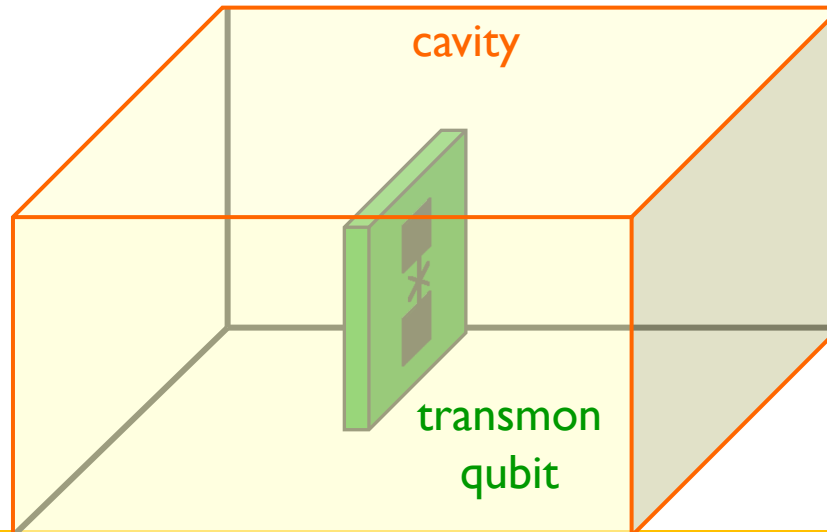
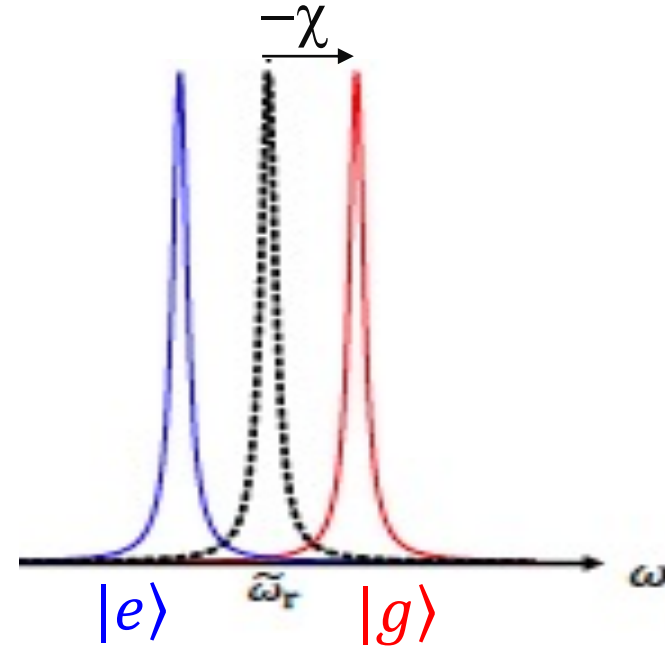
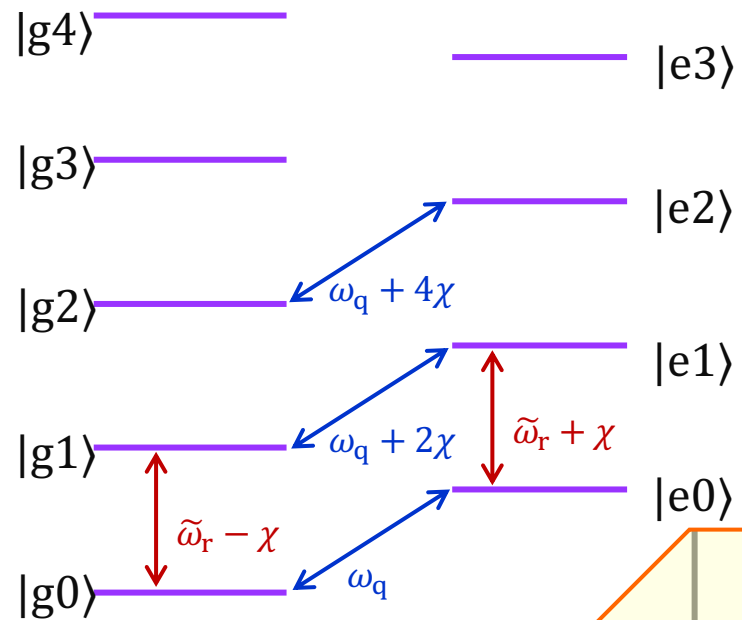
$$A = \frac{2d_0^2}{3} \left(\frac{1}{\Delta_0^2}\right) \left[\frac{r_L^2}{V_L^5} + \frac{r_T^2}{V_T^5}\right]^{-1} * \frac{2\pi\hbar^2}{\varepsilon} * \frac{\pi\hbar\rho^2 V^2}{C\gamma P k \Delta} \quad \mathbf{E^2}$$



Qubit Readout

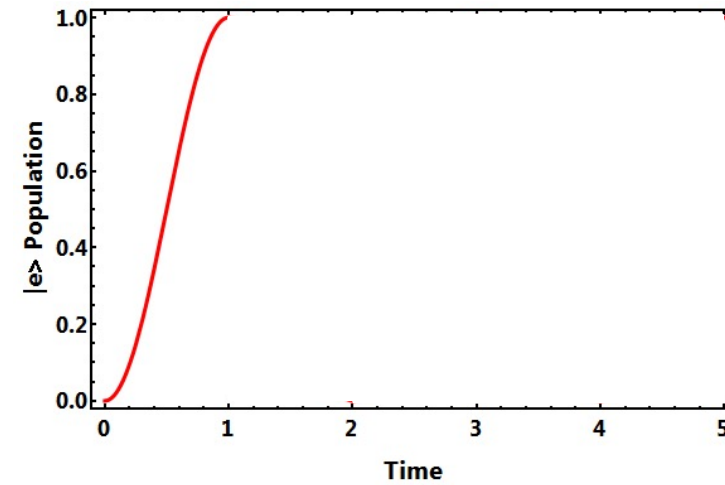
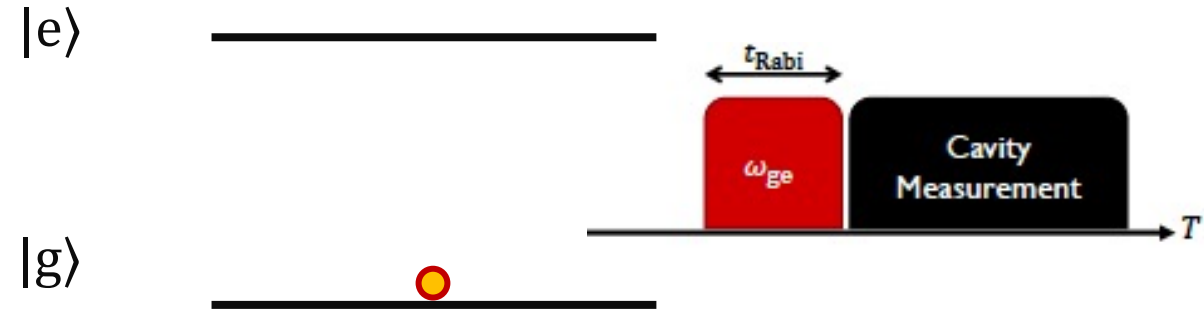
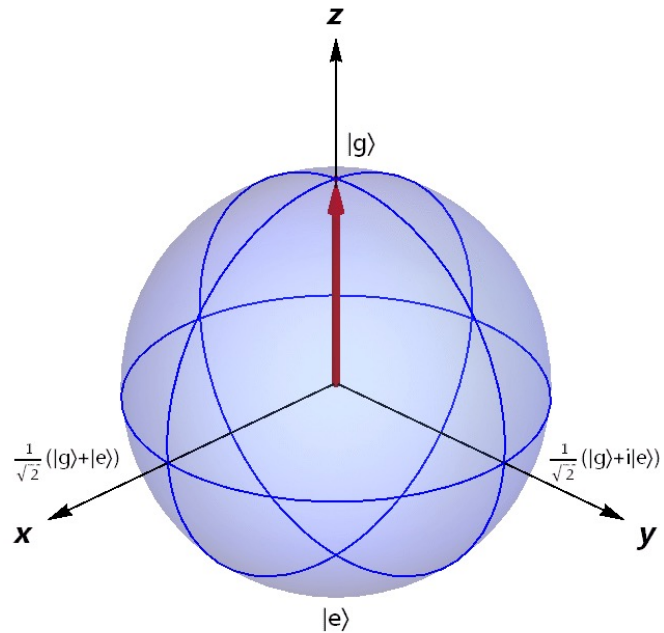
Dispersive Jaynes-Cummings Hamiltonian

$$\mathcal{H}_{\text{JC}}^{(\text{disp})} = \hbar(\tilde{\omega}_r + \chi\sigma_z)a^\dagger a + \frac{1}{2}\hbar\tilde{\omega}_q\sigma_z$$



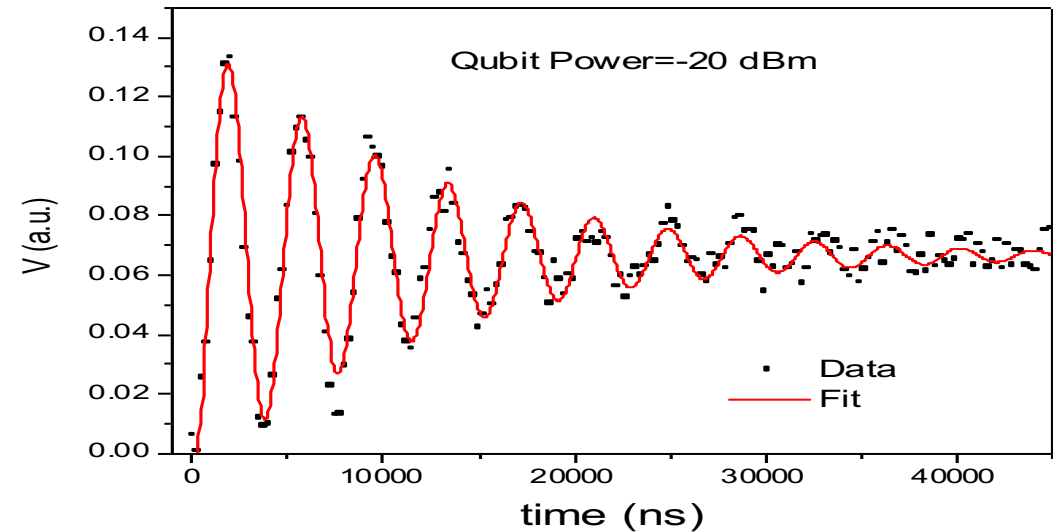
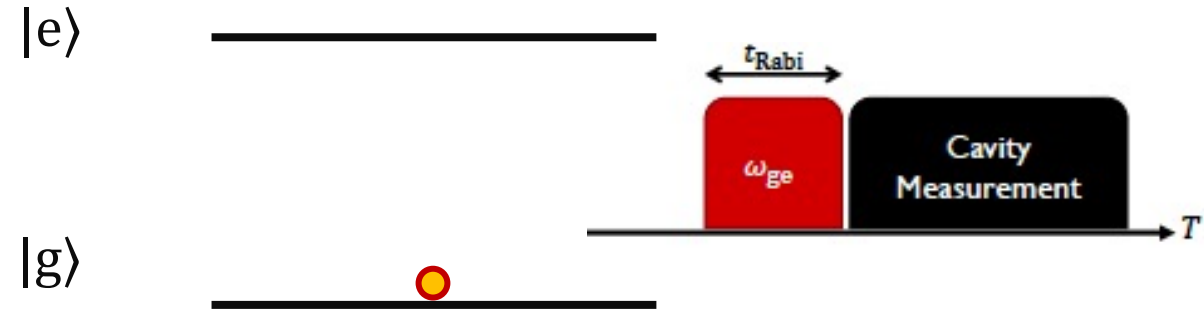
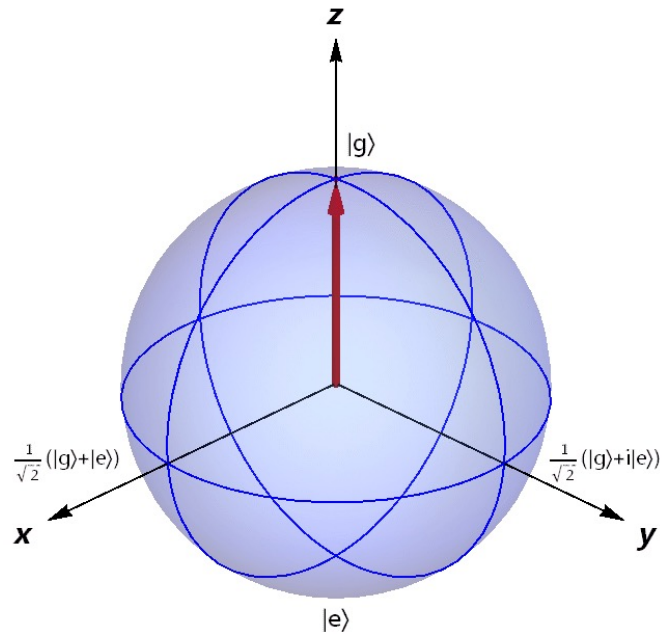
Driving the Qubit & Rabi Oscillations

Driving a qubit on resonance
at frequency ω_{ge}



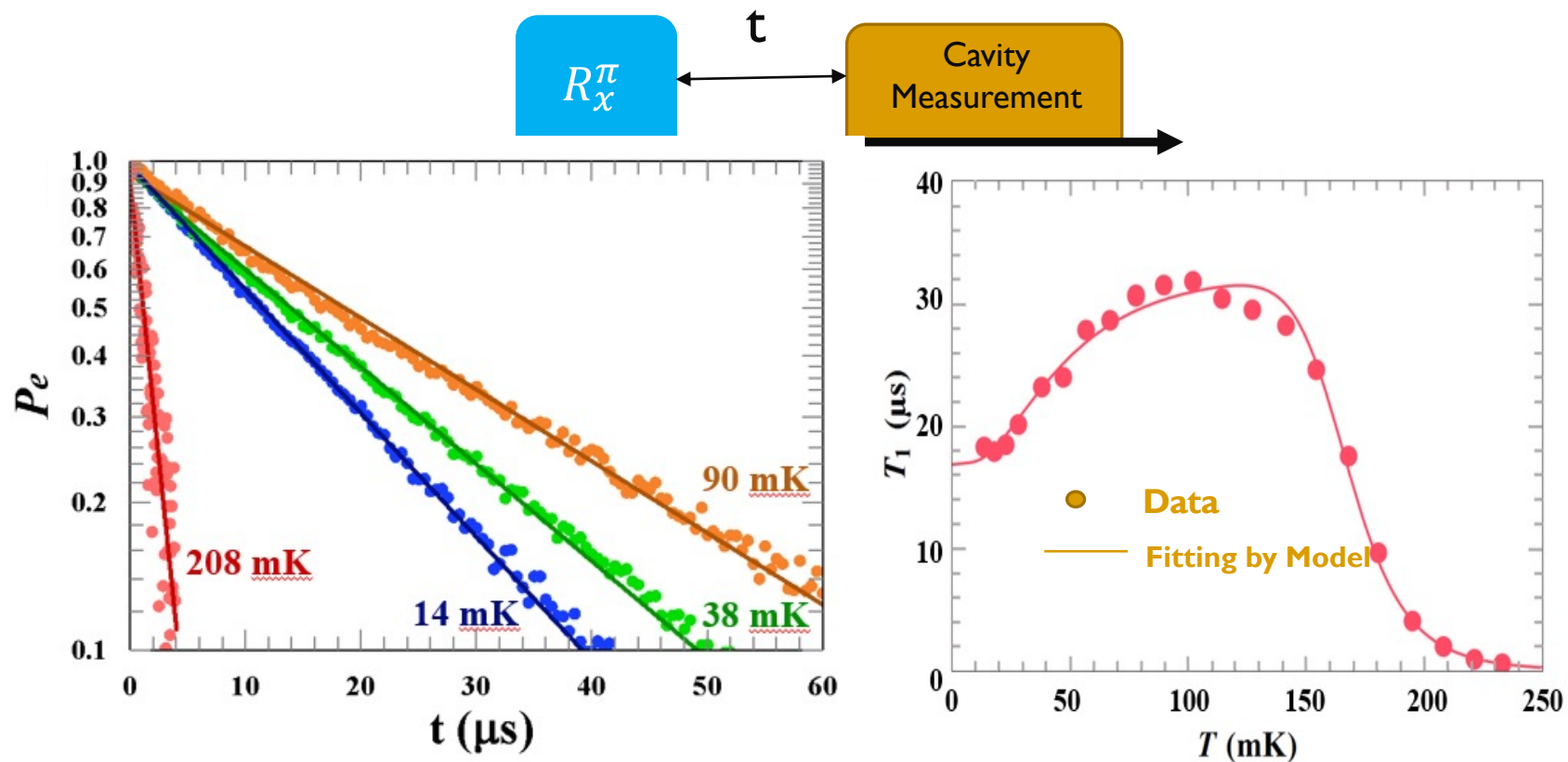
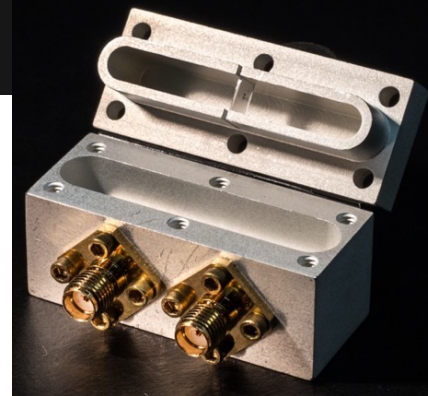
Driving the Qubit & Rabi Oscillations

Driving a qubit on resonance
at frequency ω_{ge}



Transmon in 3D Cylindrical Cavities

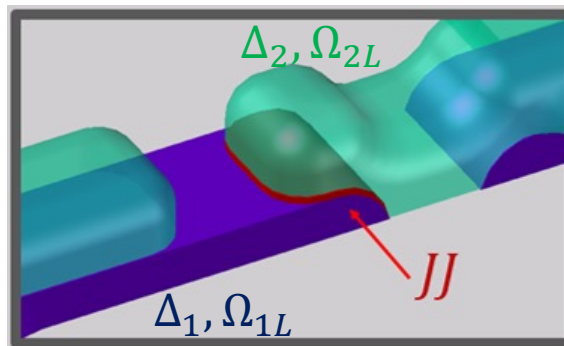
Qubit Relaxation Measurement



1. TLS Loss?

2. Imbalanced Nonequilibrium Quasiparticles at two different SC gaps?

$$n_{1L} \approx n_{th,1} + \frac{n_{ne,1L}}{1 + \frac{\Omega_{2L}}{\Omega_{1L}} \sqrt{\frac{\Delta_2}{\Delta_1}} \exp\left(\frac{\Delta_1 - \Delta_2}{K_B T}\right)}$$

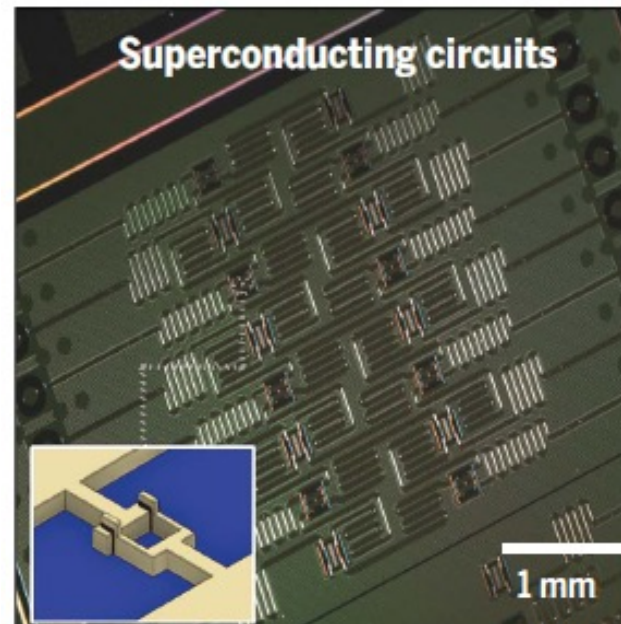


$$n_{2R} \approx n_{th,2} + \frac{n_{ne,1R}}{\sqrt{\frac{\Delta_1}{\Delta_2}} \exp\left(\frac{\Delta_2 - \Delta_1}{K_B T}\right) + \frac{\Omega_{2R}}{\Omega_{1R}}}$$

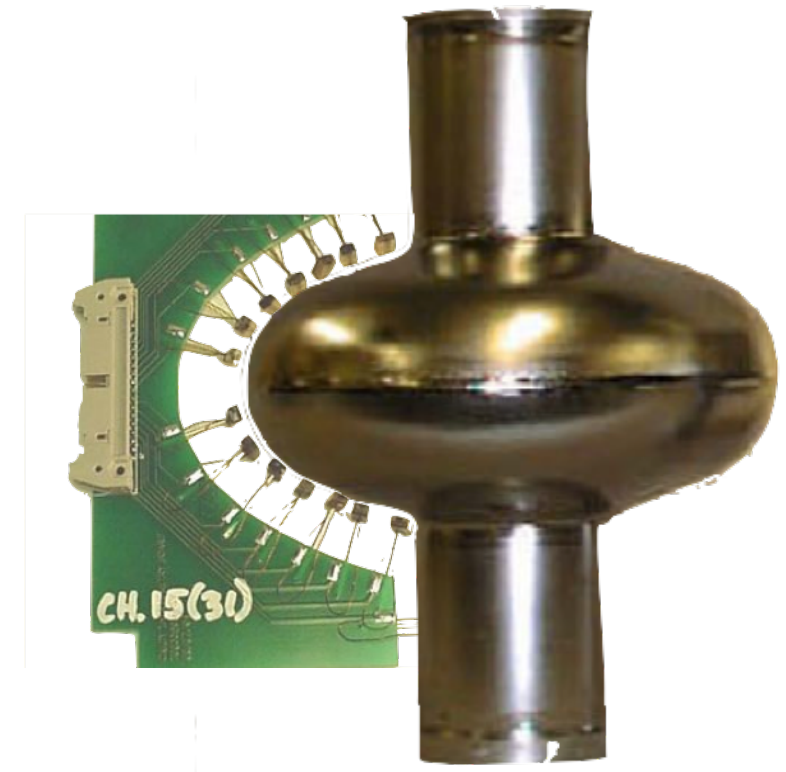
Motivation of Microwave Microscope

$$\frac{1}{Q_i} = \frac{1}{Q_{TLS}(T,E)} + \frac{1}{Q_{QP}(T,E)} + \frac{1}{Q_{rad}(W,geometry)} + \frac{1}{Q_{vortices}(B)} + \dots$$

Develop a microscopy which can microscopically locate the place of loss at devices/cavities operating condition and identify the loss mechanism.

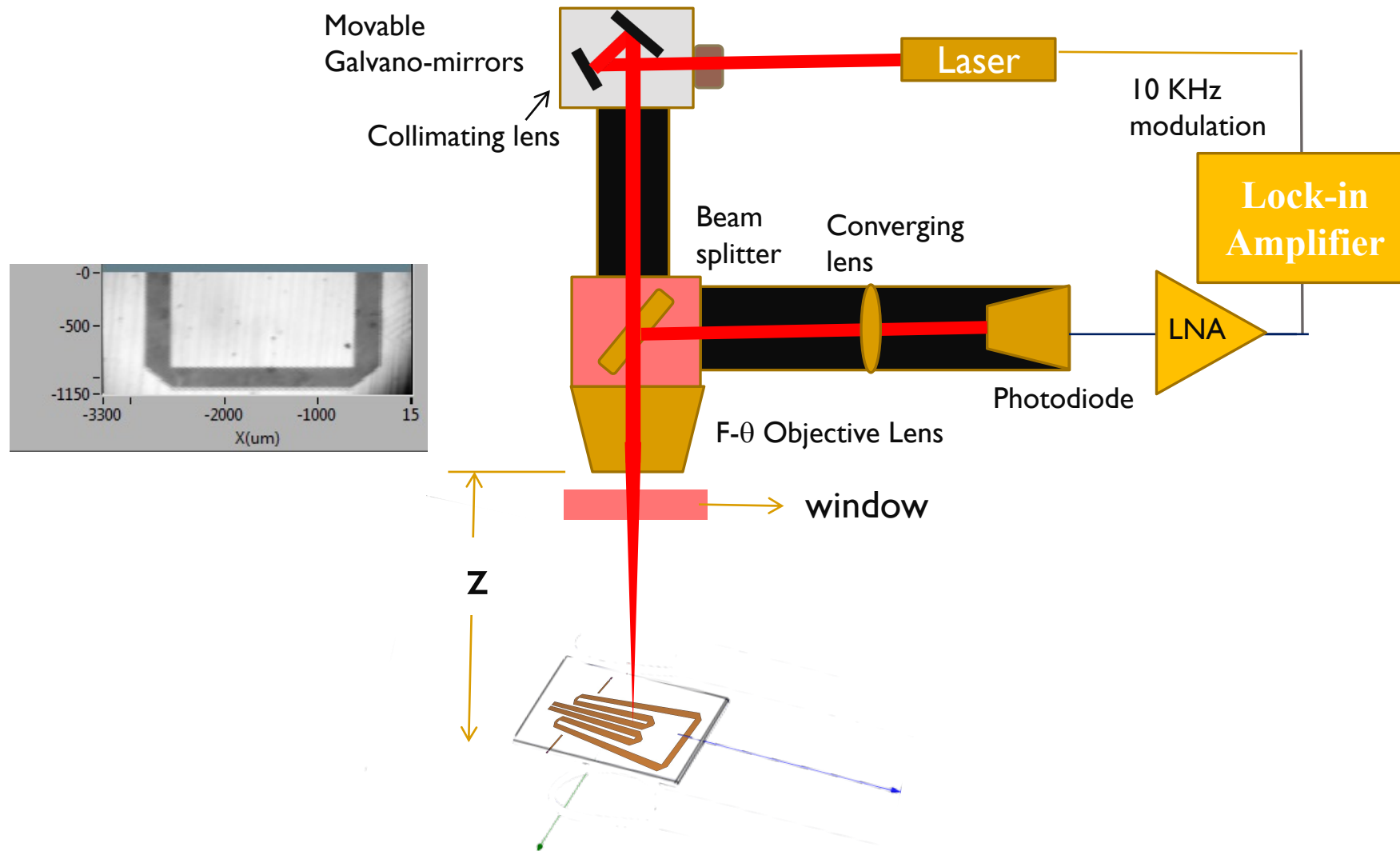


de Leon *et al.*, *Science* **372**, eabb2823 (2021) 16 April 2021



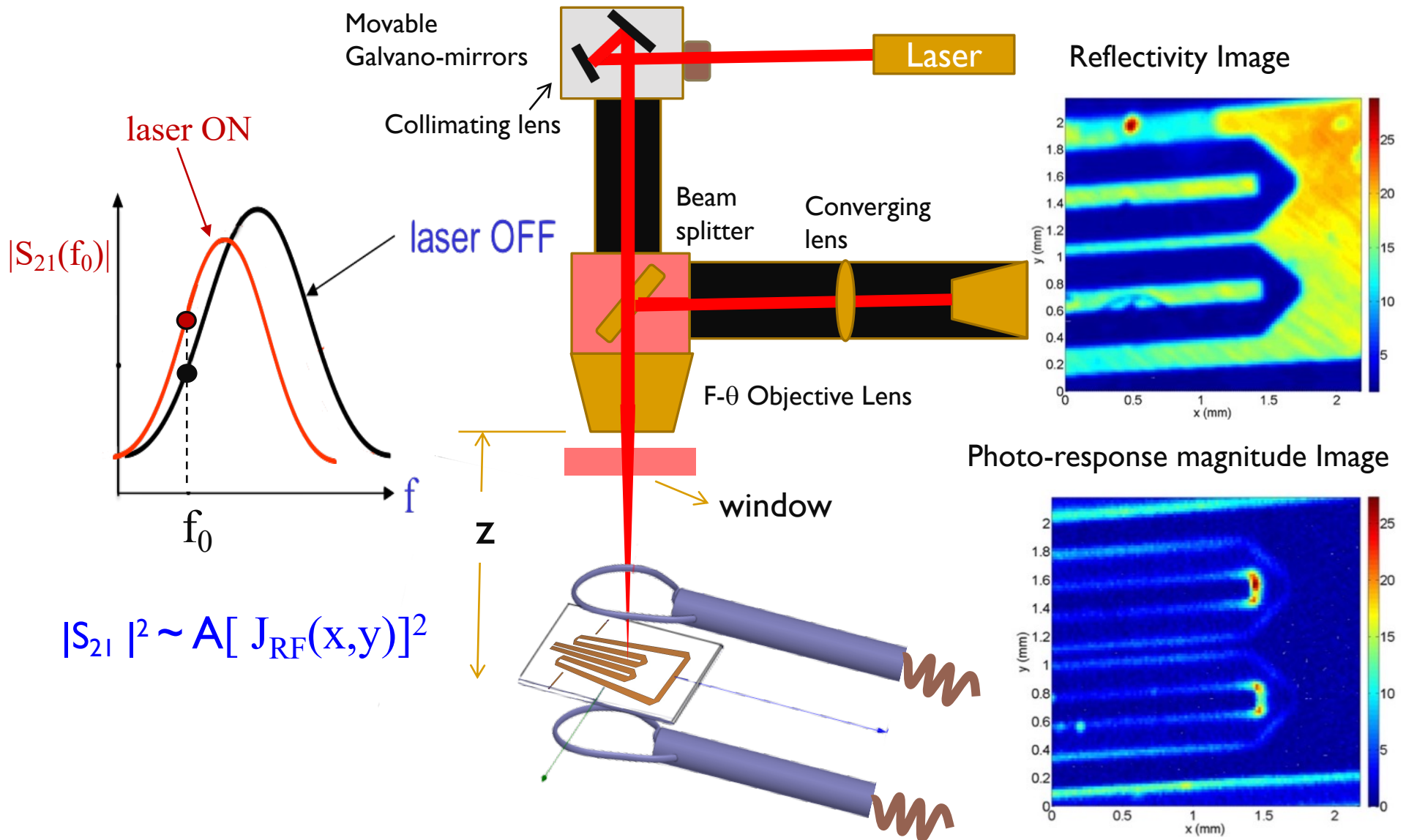
Laser Scanning Microwave Microscopy

1. Reflectance mode to generate a conventional image

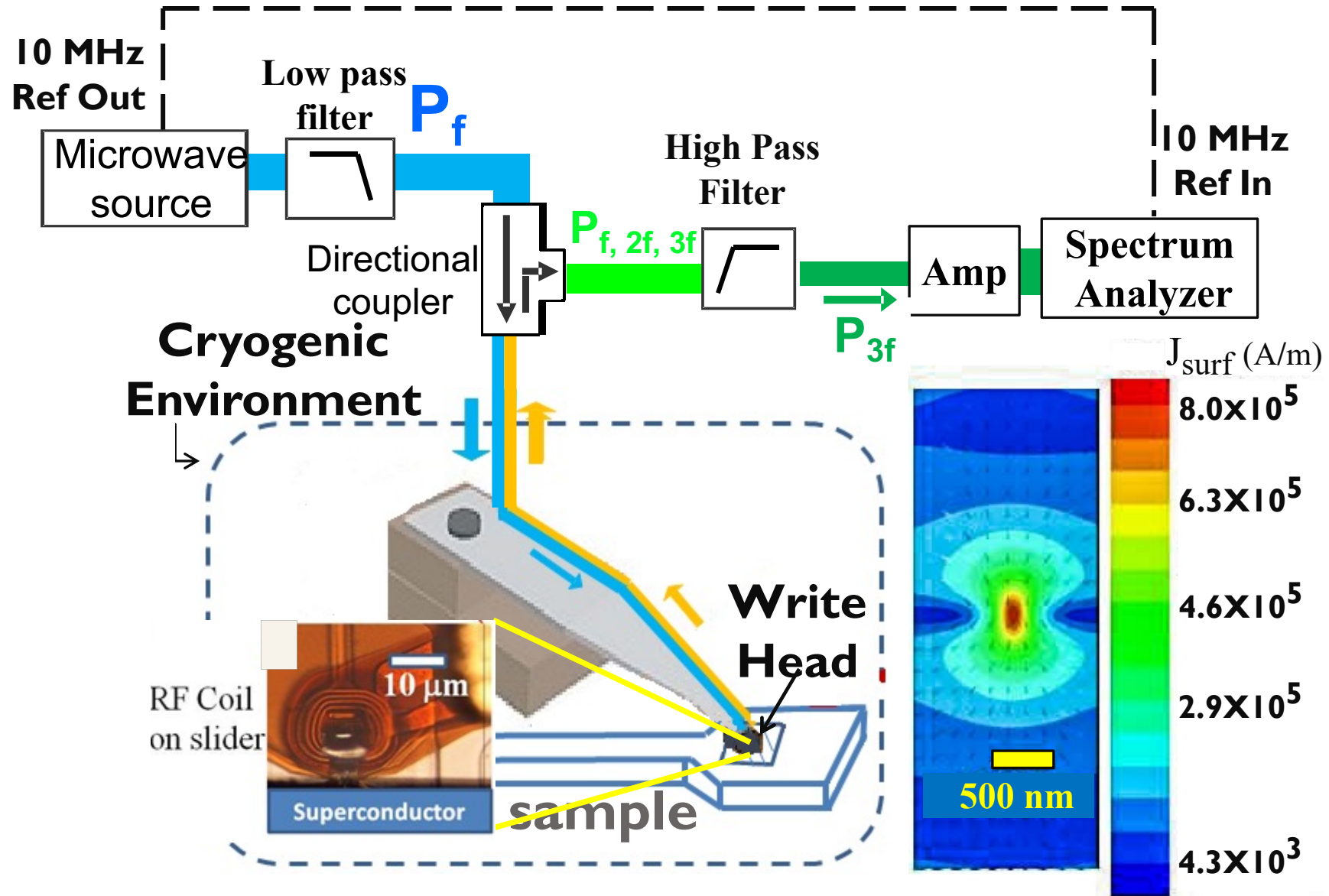


Laser Scanning Microwave Microscopy

2. Photoresponse mode to image the current density

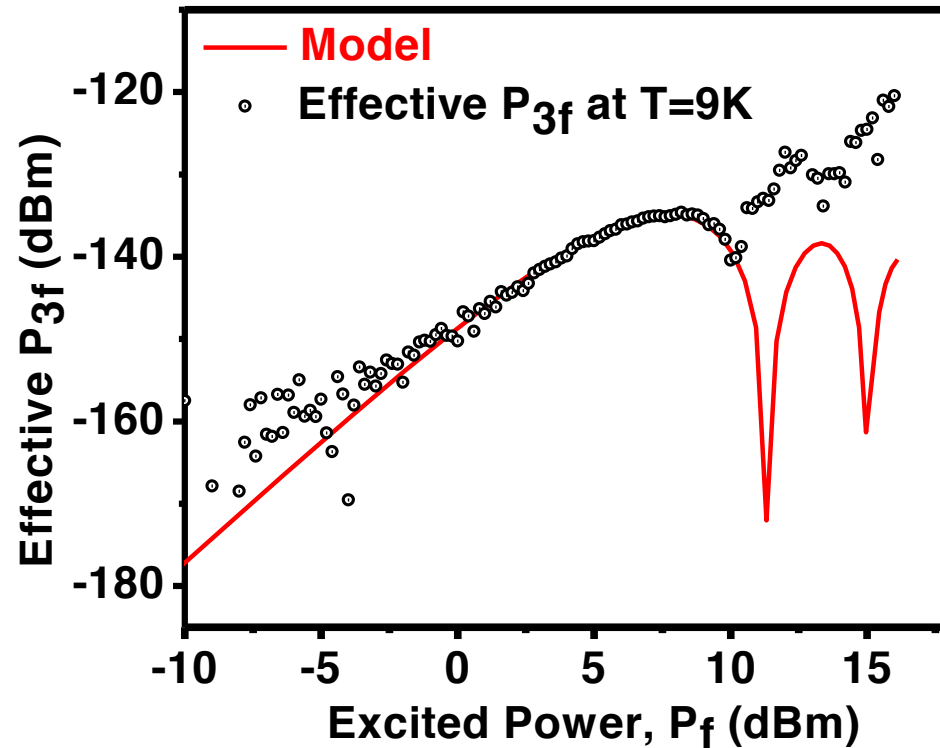


Near-Field Microwave Microscopy

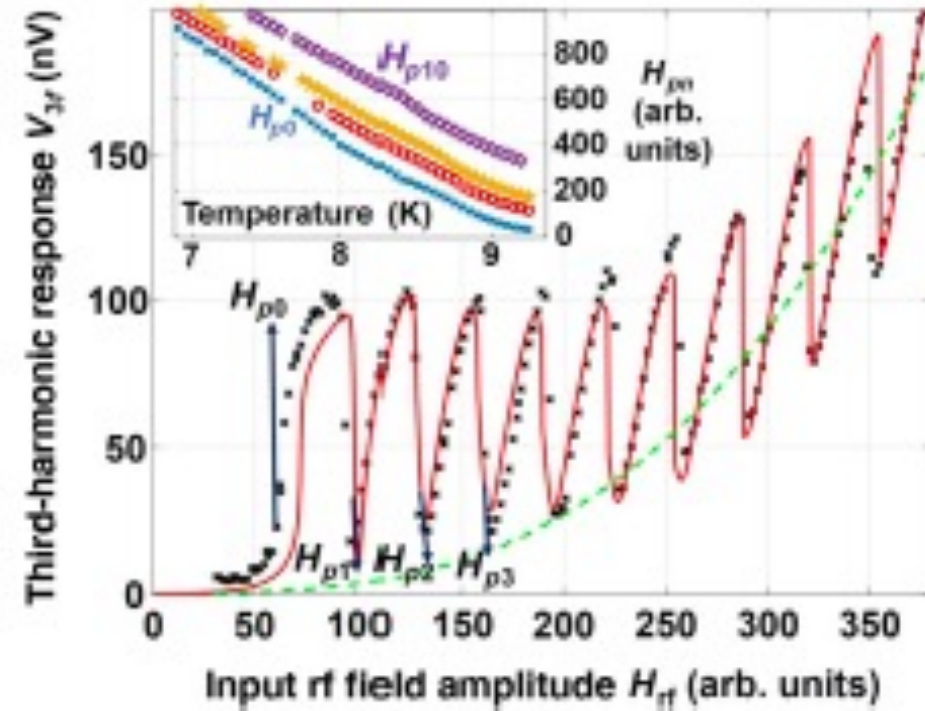


Nonlinear Measurement and Modeling

General curve of $P_{3f}(P_f)$



Tamin Tai et al, PRB 92, 134513 (2015)



Bakhron Oripov et al.

PHYSICAL REVIEW APPLIED 11, 064030 (2019)

Conclusions

❑ Superconducting Quantum Computing

- ❖ In aluminum resonator, an unusual increase of Q_i with decreasing temperature is observed, which is due to the increase of TLS coherence time (T_2) in ultra-low temperature and power.
- ❖ Temperature dependent of T_1 on transmon qubits is also affected by many loss issues. In addition to TLS loss issue, two gaps model also can interpret the T_1 temperature dependent behavior.

❑ Microscopy

❖ Laser Scanning Microscopy

A clear photo response with the capability of phase-sensitive measurement of local microwave properties on YBCO resonators are obtained by LSM with $\sim \mu\text{m}$ length resolution.

❖ Near-field Scanning Microwave Microscopy

This microscopy achieved local harmonic generation from bulk Nb surfaces from sub-micron scale length and the measured nonlinearity is interpreted by the model of weak-link Josephson junctions.