



Nuclear Structure and Reactions at NSCL and FRIB through the Lens of Astrophysics: Lecture 1

Chris Wrede

HUGS @ JLab

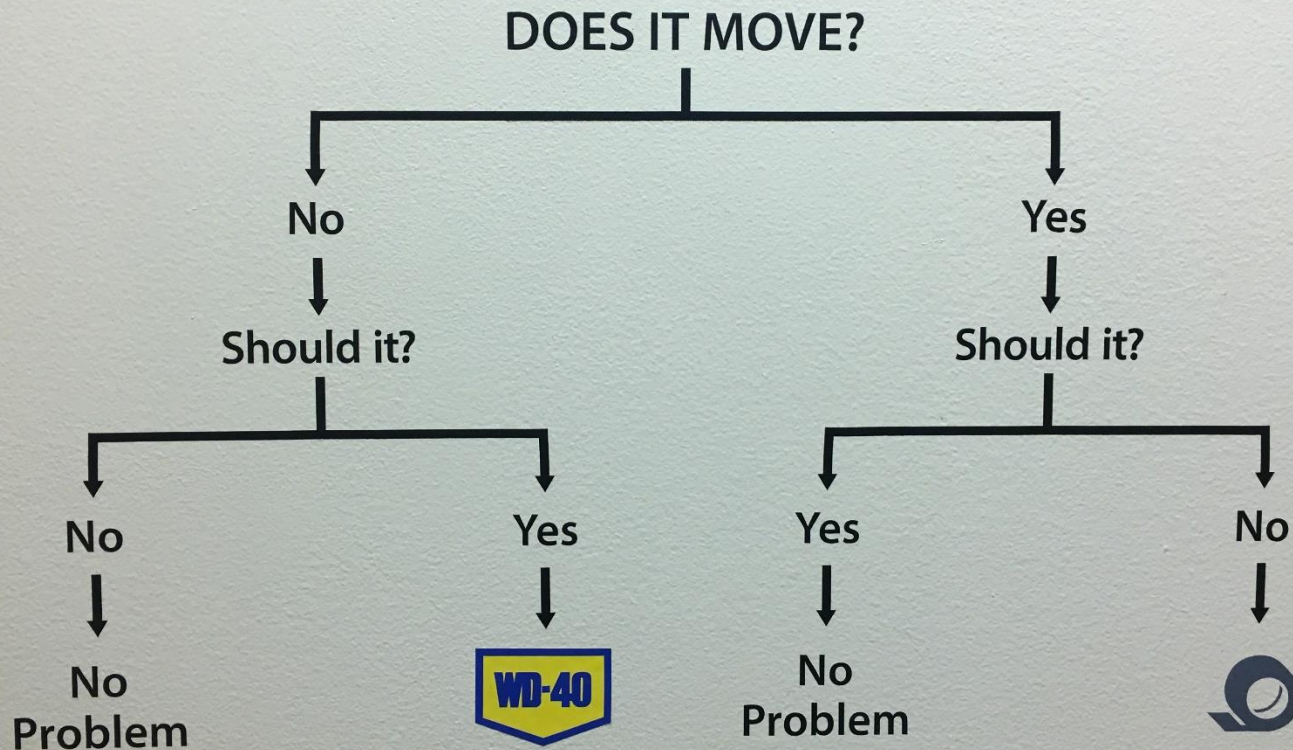
May 29th, 2019



MICHIGAN STATE
UNIVERSITY

Painted on the wall of my room at SURA

Engineering Flowchart

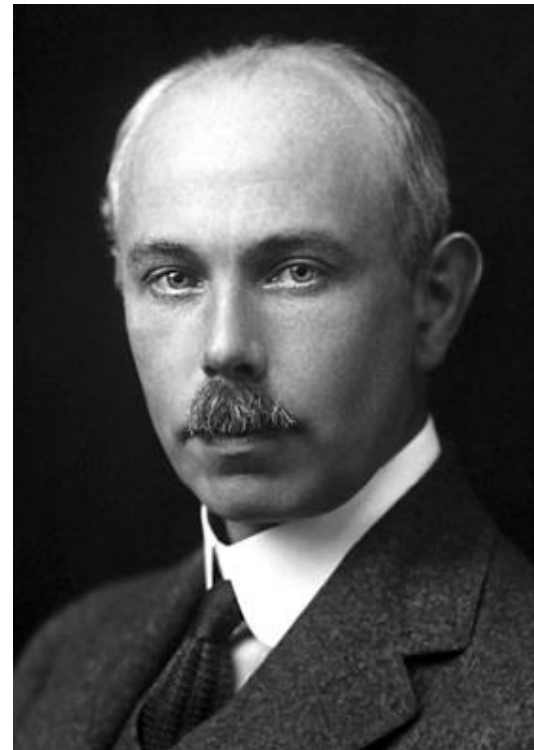
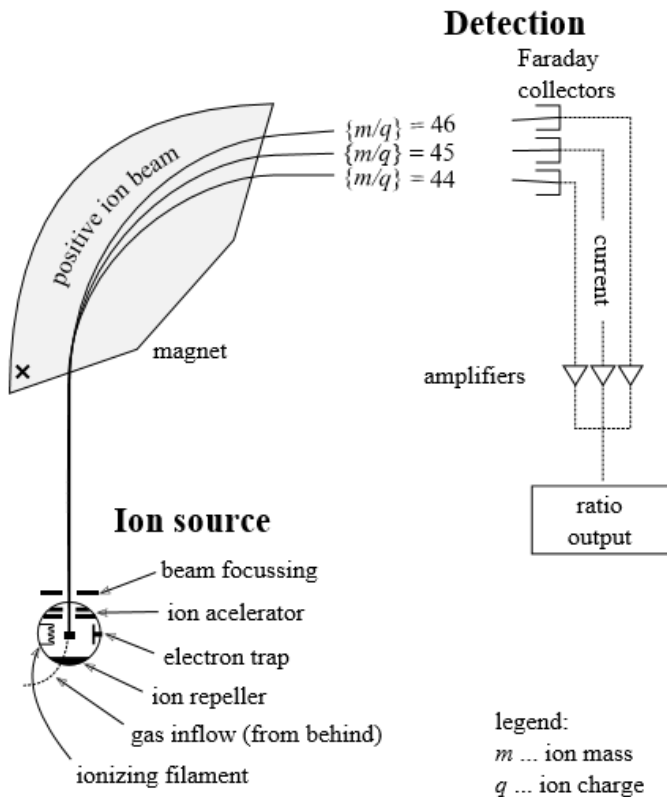


Outline

- **Lecture 1: History, stellar evolution & thermonuclear rates**
- Lecture 2: Charged-particle reactions: direct measurements
- Lecture 3: Charged-particle reactions: indirect measurements
- Lecture 4: Slow neutron capture process: direct measurements
- Lecture 5: Rapid neutron capture process: indirect measurements

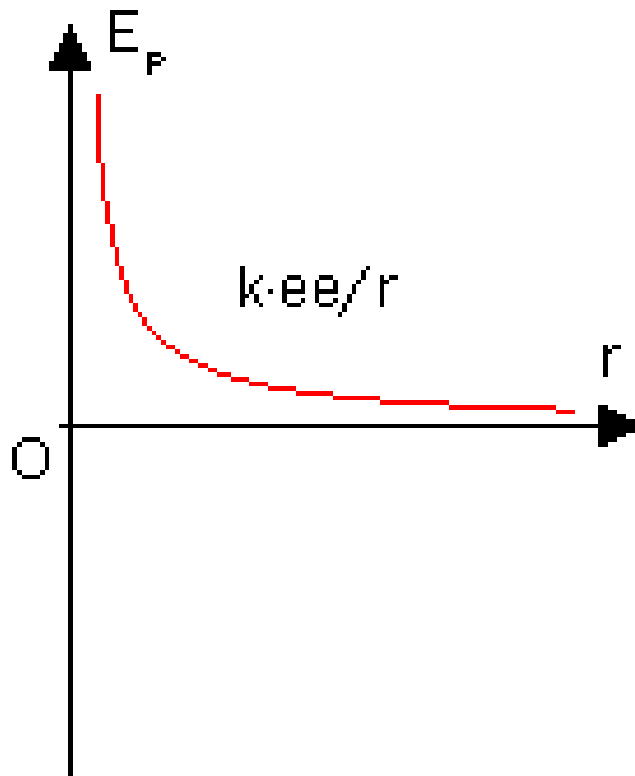
H, He atomic masses

- Aston (1920): Uses mass spectrometry to reveal that He atomic mass slightly less than 4 times H mass



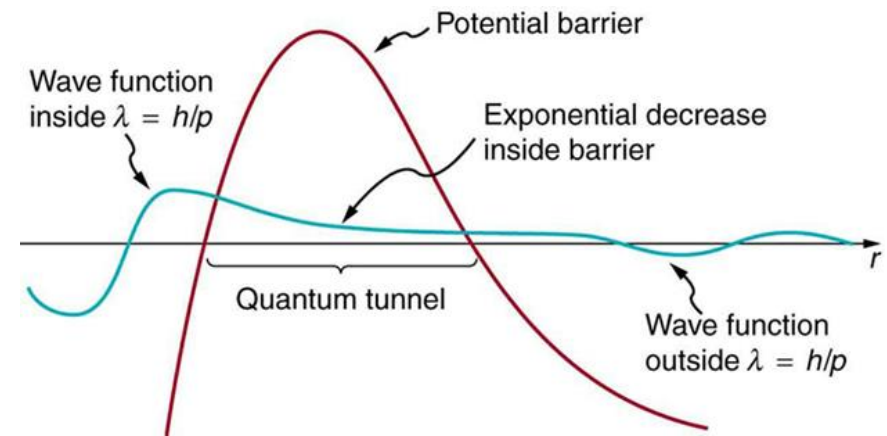
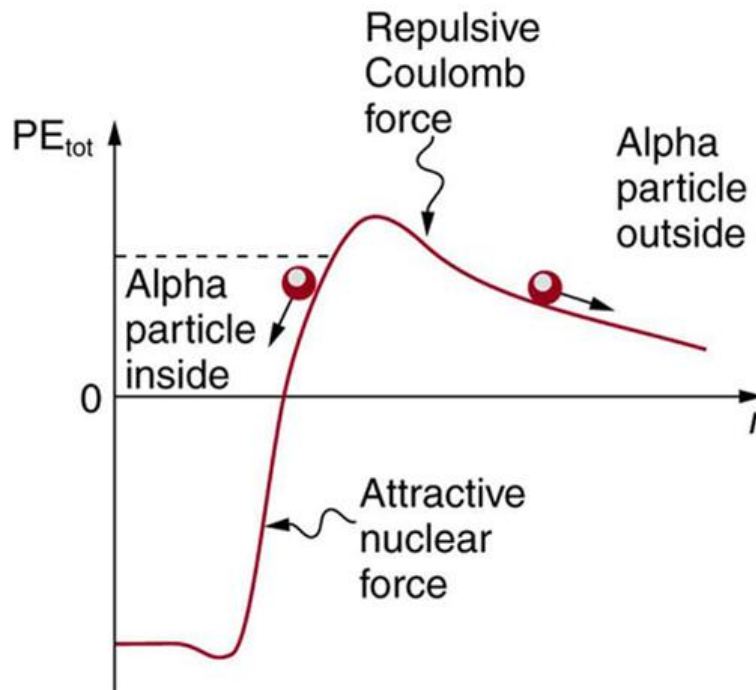
Solar energy: $E = mc^2$

- Eddington (1920): conversion of H to He will explain energy generation in Sun!
- Problem: how do protons overcome Coulomb repulsion?



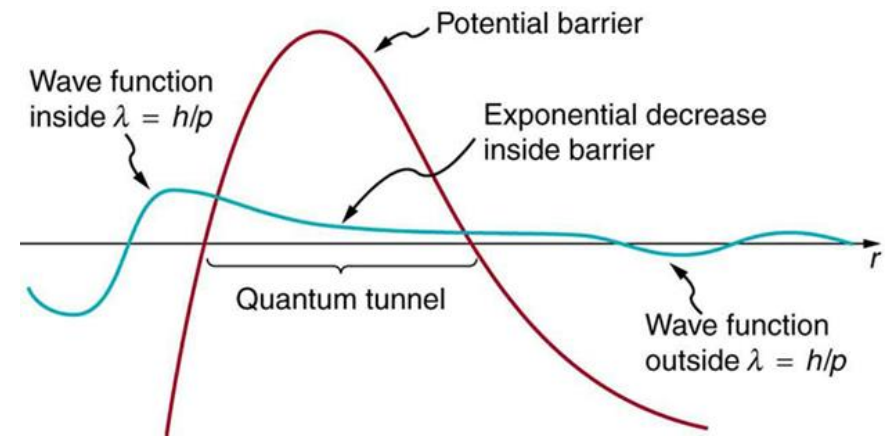
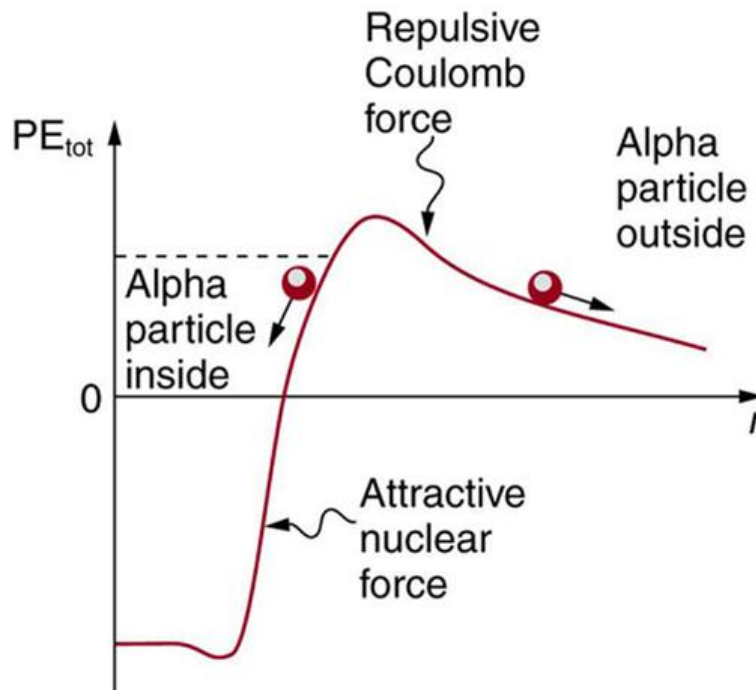
QM tunneling: decay

- Gamow (1928) & Condon and Gurney (1929): explained alpha decay via quantum mechanical tunneling



QM tunneling: fusion

- Atkinson and Houtermans (1929): tunneling through Coulomb barrier may enable nuclear fusion in stars



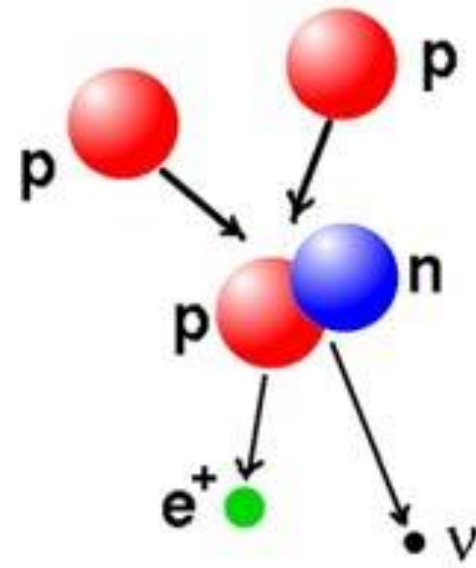
Stellar reactions in the lab

- Cockroft and Walton (1932): first nuclear reaction using artificially accelerated particles
- break-up of Li into two alpha particles under bombardment by ~few 100 keV protons
- part of stellar hydrogen-burning *pp chains*



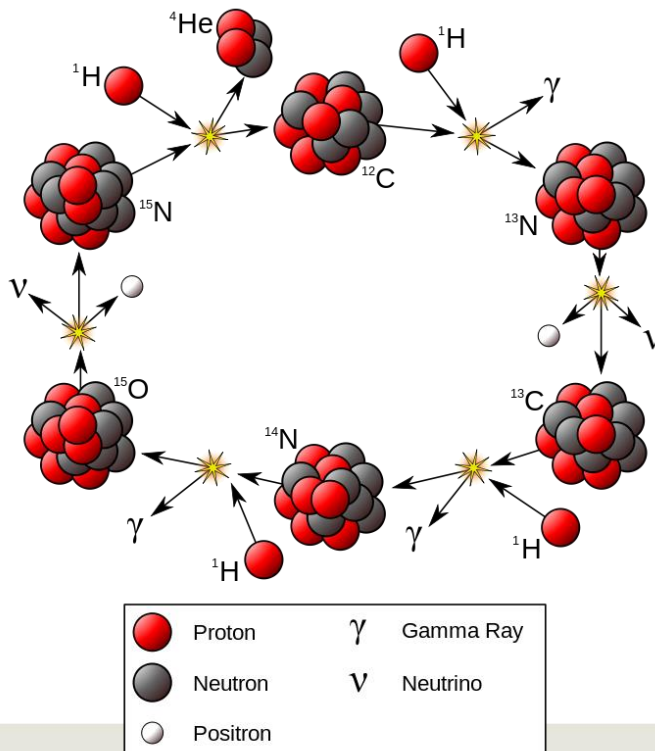
pp chains: the first step

- Atkinson (1936) proposed fusion of two hydrogen nuclei to deuterium as a source of stellar energy generation
- Bethe and Critchfield (1938) showed this reaction gives energy generation of correct order of magnitude to power Sun
- First step of stellar hydrogen-burning *pp chains*



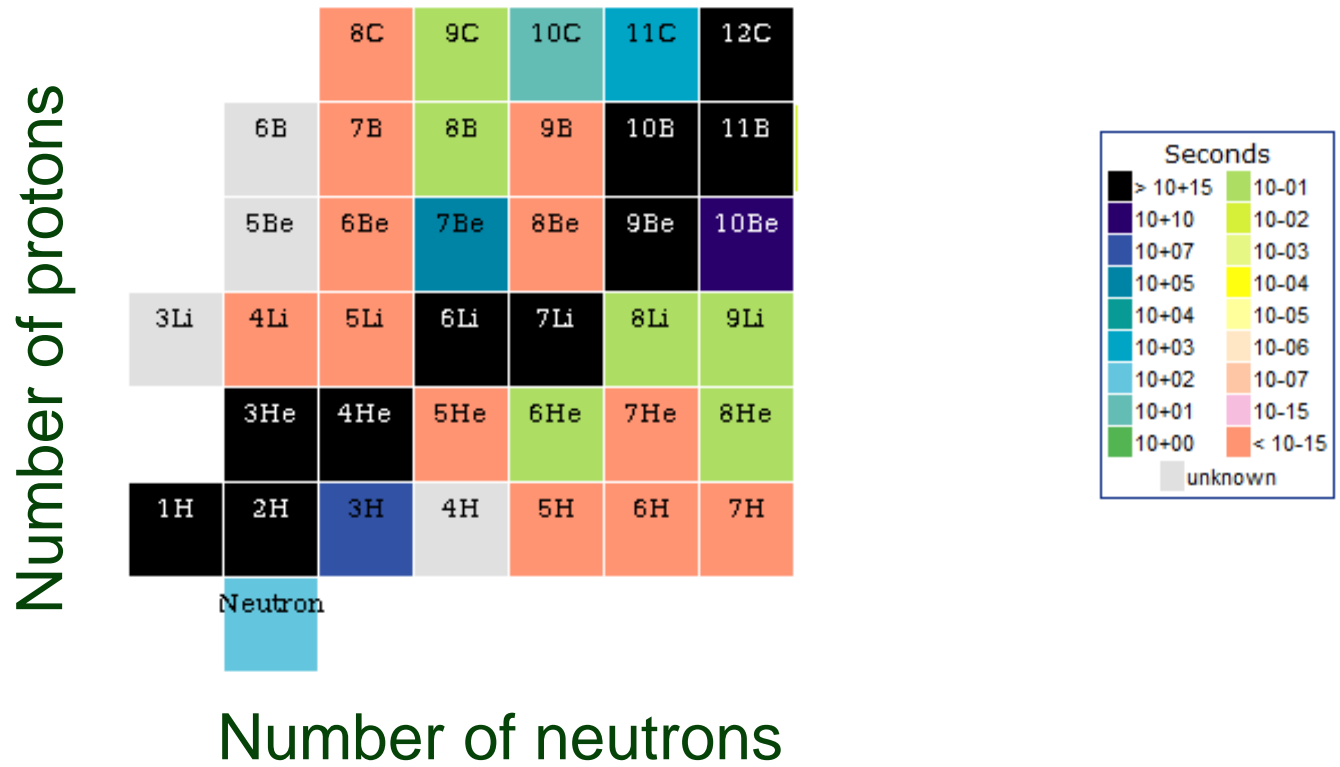
CNO cycle

- Lauritsen and Crane (1934) produced 10-min radioactivity (^{13}N) by bombarding carbon with protons
- Bethe (1939) and von Weizsacker (1938): energy production in stars by CNO cycles discovered: rate and temperature dependence



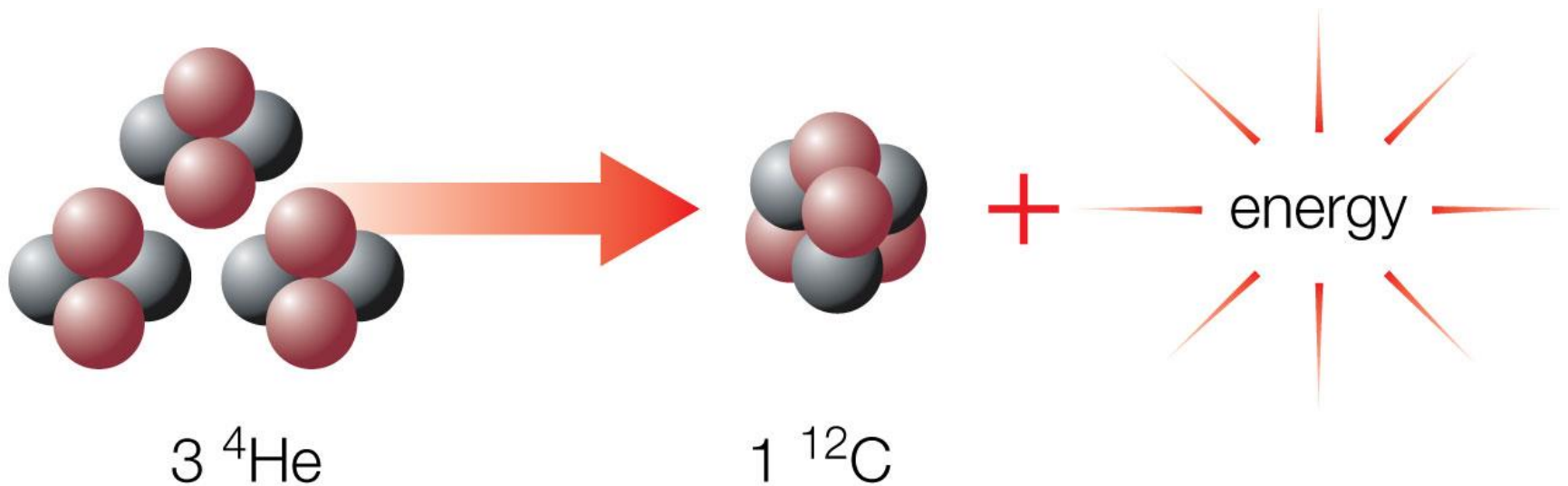
A = 5, 8 nucleosynthesis gaps

- Hoyle (1946, 1954): theory of nucleosynthesis within framework of stellar evolution using available nuclear data
- Problem: no stable nuclei of mass number 5, 8; how to bypass these in stars?



3α reaction

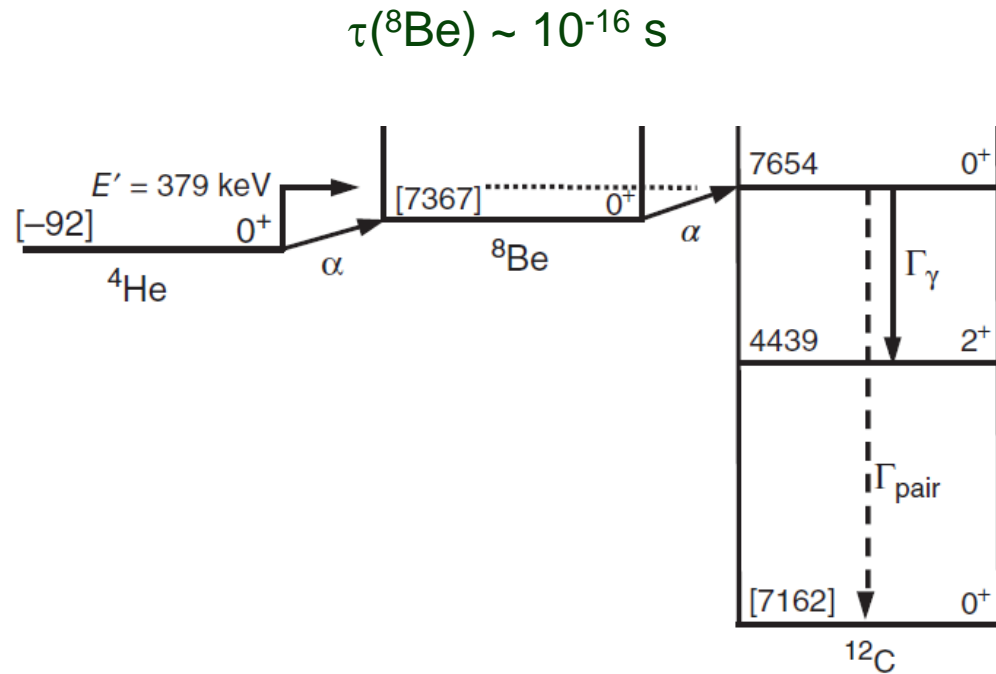
- Salpeter (1952): three helium atoms could combine to form carbon via “triple alpha reaction” powering red giant stars



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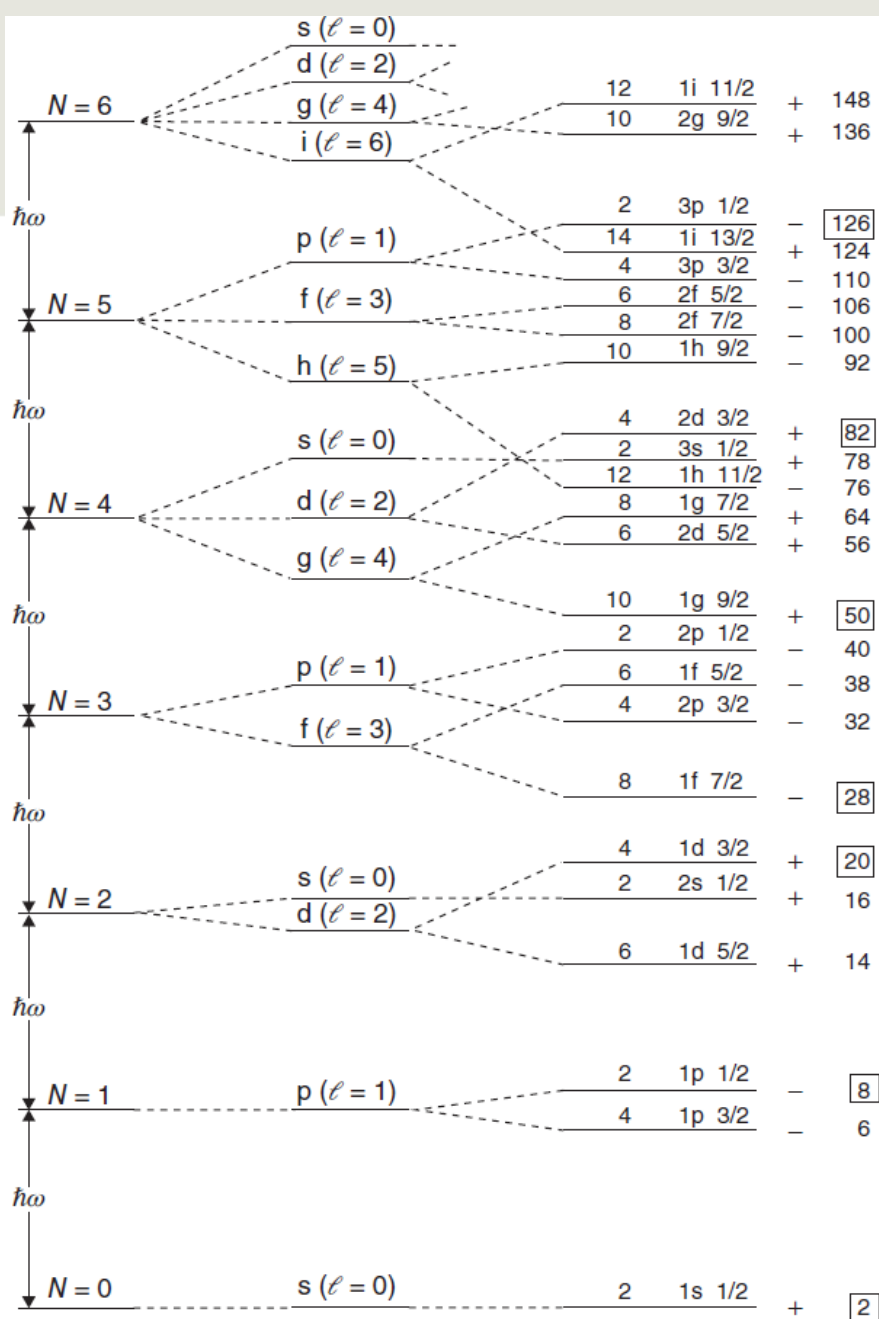
3α reaction: Hoyle state

- Hoyle: must be special ^{12}C excited resonance state at 7.7 MeV to make triple alpha reaction fast enough to power red giant stars
- Dunbar *et al.* (1953), Cook *et al.* (1957): discovered the “Hoyle State” experimentally confirming that the triple alpha reaction can bypass mass 5, 8 gaps



Nuclear shell model

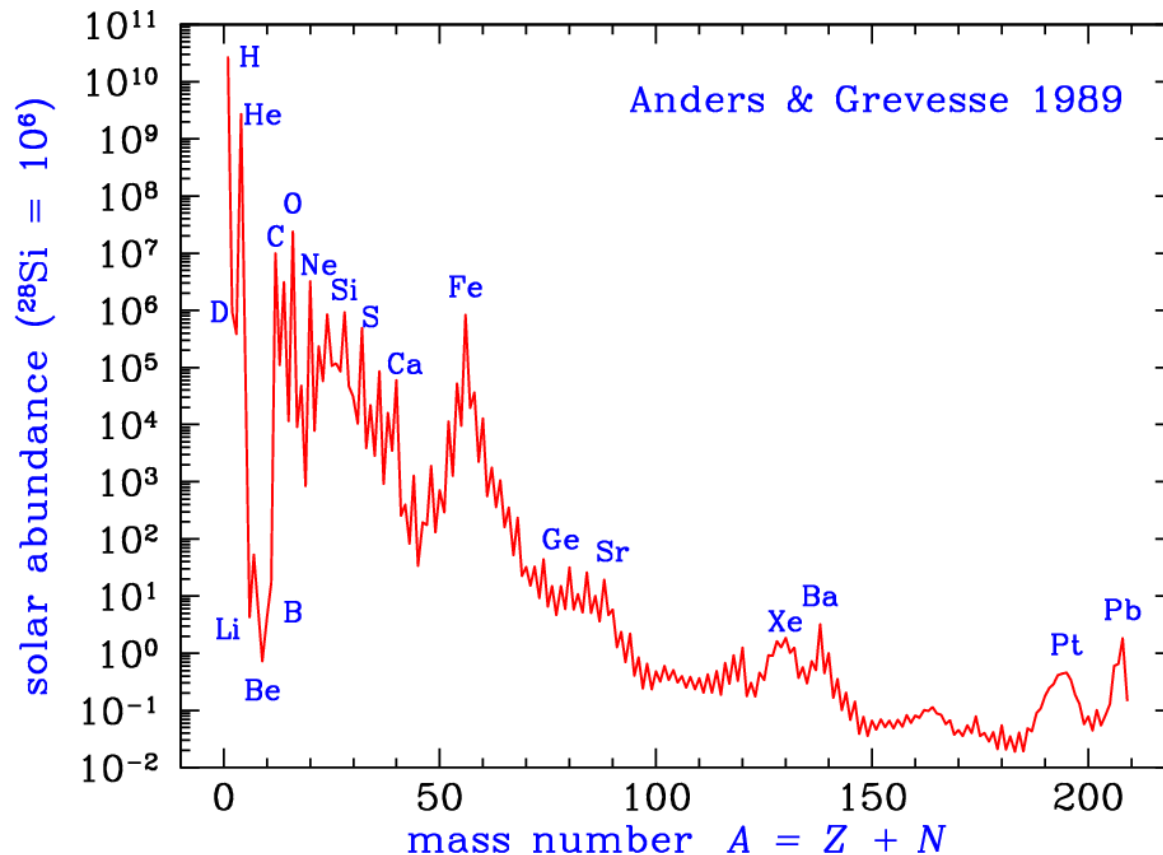
- Jensen and Goeppert Mayer (1949) develop nuclear shell model and discover magic numbers



Harmonic oscillator $\ell = N, N-2, \dots$ $\vec{\ell} \cdot \vec{s}$ $N_j = 2j+1$ $n\ell j$ $\Pi = (-1)^\ell \Sigma N_j$

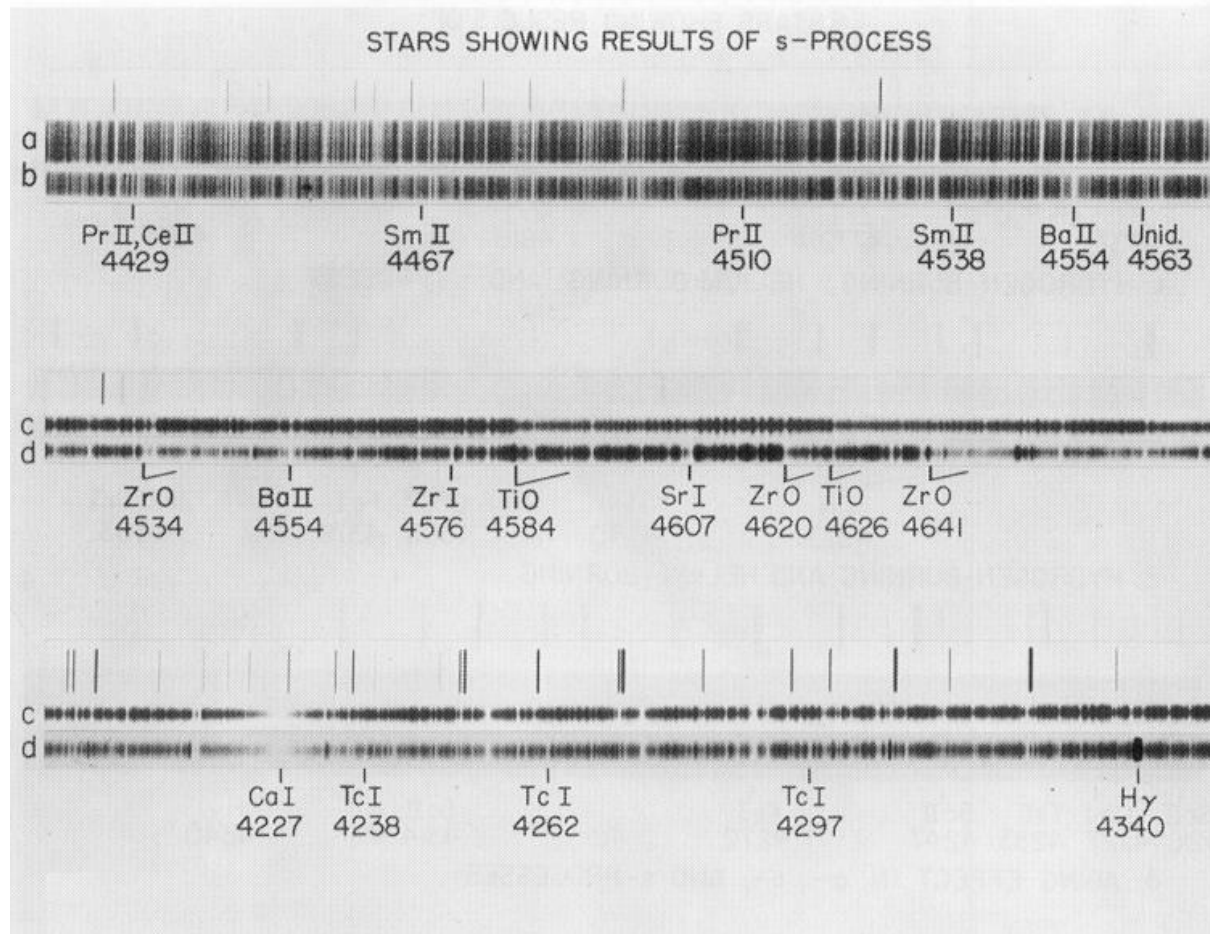
Solar abundances & nuclear shell model

- Suess and Urey (1956): discovered peaks in solar system abundances, related to magic numbers motivating nucleosynthesis theory for heavy elements via *s*, *r* processes



Stellar nucleosynthesis: a smoking gun

- Merrill (1952): discovery of unstable Tc in red giant stars showing that nuclear reactions are producing heavy elements in stars



Foundations of nuclear astrophysics: B²FH

REVIEWS OF MODERN PHYSICS

VOLUME 29, NUMBER 4

OCTOBER, 1957

Synthesis of the Elements in Stars*

E. MARGARET BURBIDGE, G. R. BURBIDGE, WILLIAM A. FOWLER, AND F. HOYLE



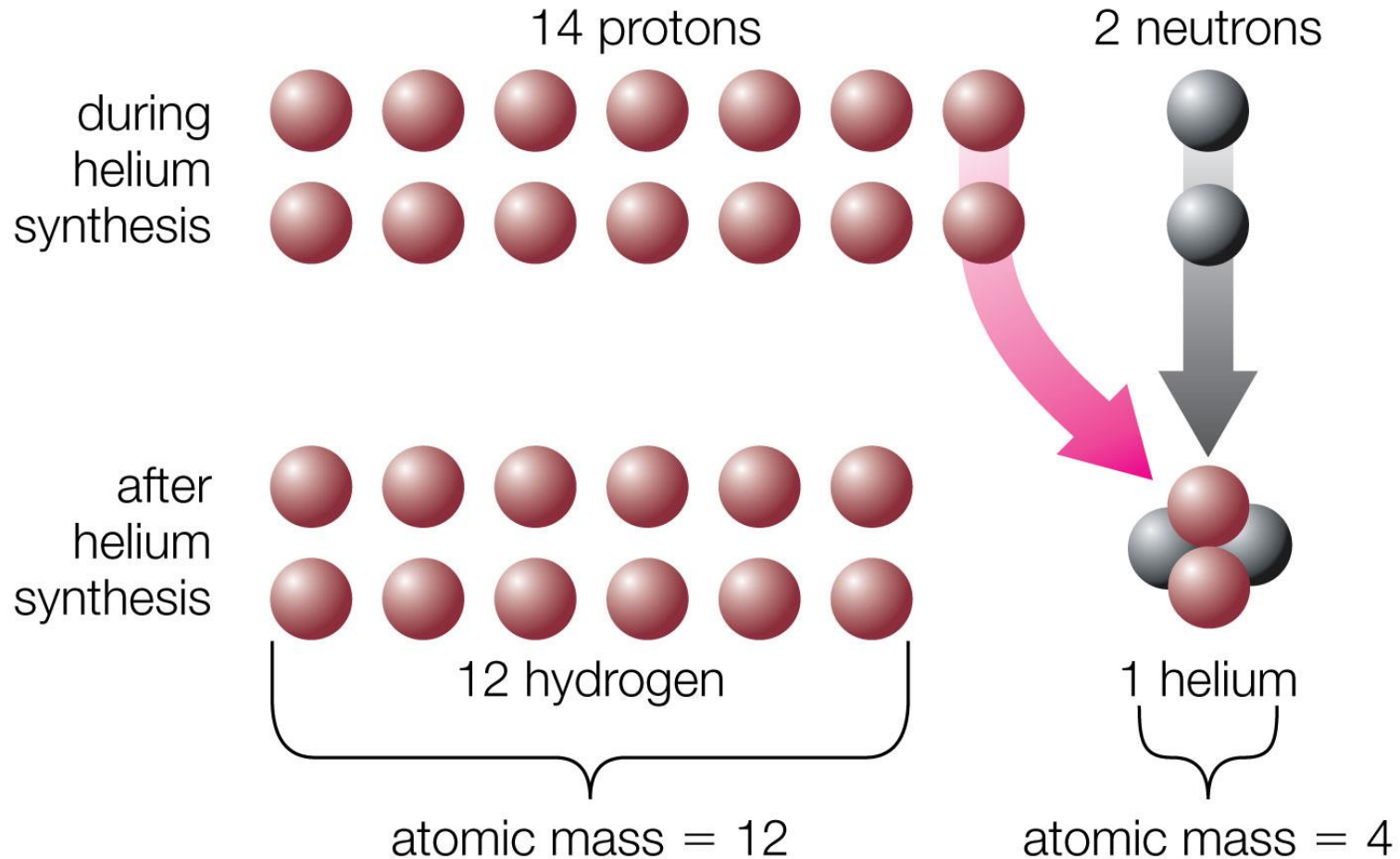
+ Al Cameron
independently
in Chalk River
internal report

Nuclear reactions in stars are responsible for energy generation and creation of elements

Nuclear astrophysics today

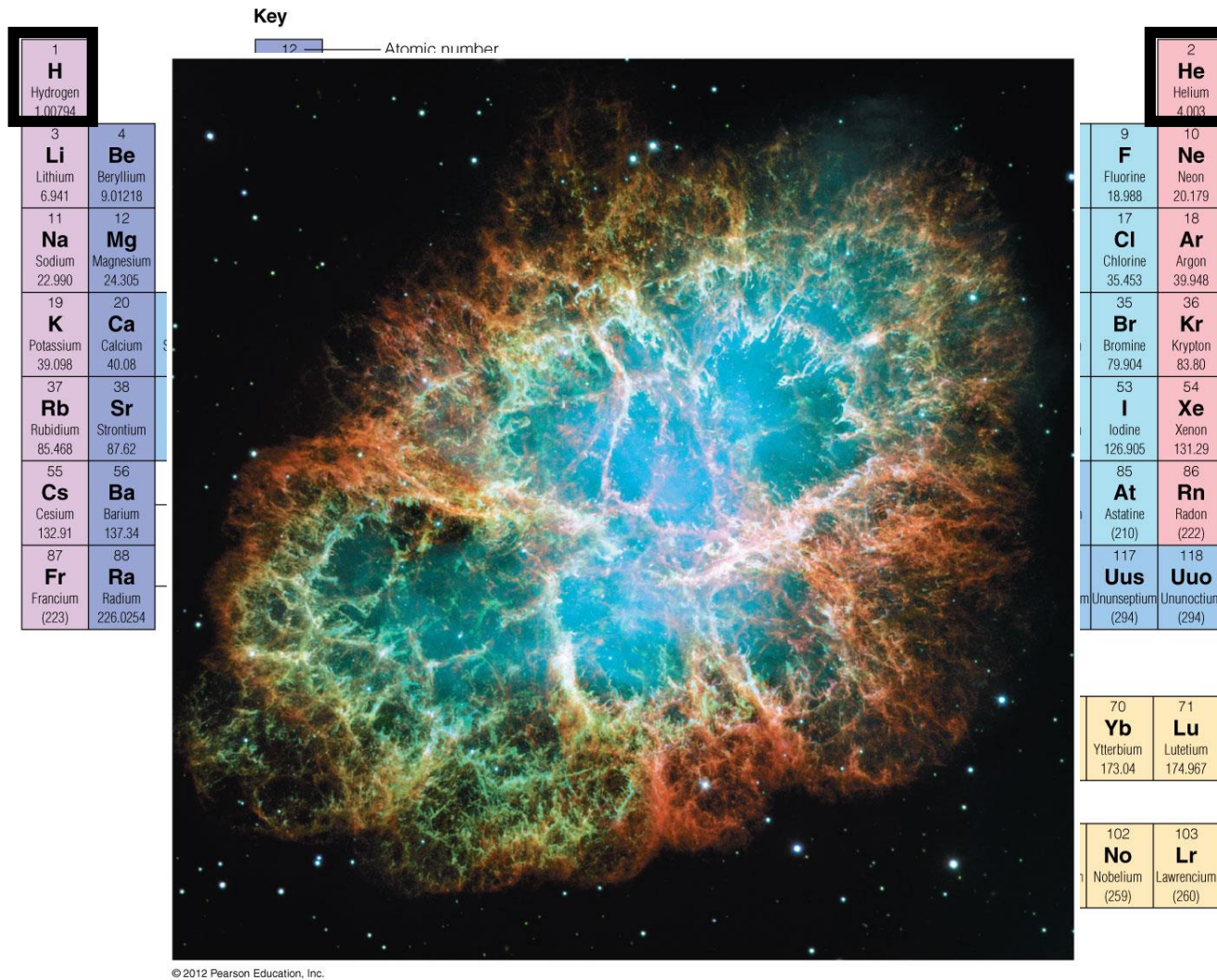
- Nuclear astrophysics has grown into a broad and vibrant field that links topics such as astronomical observation, nuclear physics experiment, nuclear theory, stellar evolution, and hydrodynamics

Big Bang nucleosynthesis



Big Bang made 75% H and 25% He, by mass, in about 1000s

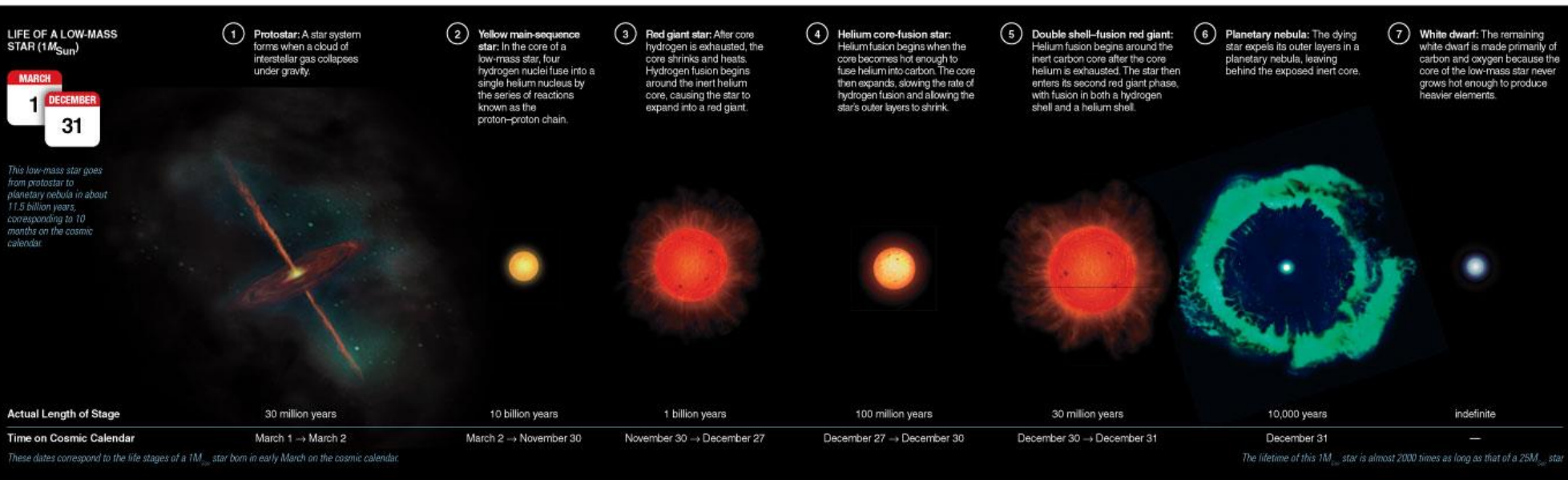
In stars and stellar events!



Life of a low-mass star

(less than about 8 solar masses)

1. Protostar: cloud of cold gas collapses under gravity
2. Main Sequence: H fuses to He in core
3. Red Giant: H fuses to He in shell around inert He core until He flash
4. Helium Core Fusion: He fuses to C, O in core while H fuses to He in shell
5. Double Shell Fusion: H and He both fuse in shells around inert C, O core
6. Planetary Nebula: outer layers expelled
7. White dwarf star left behind



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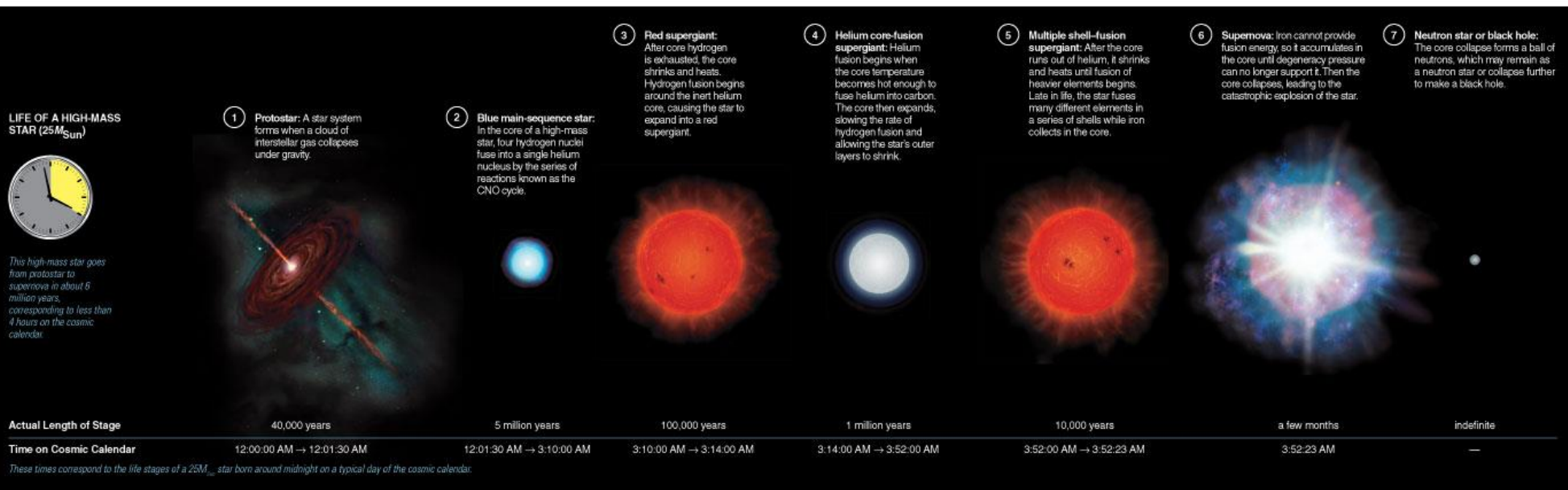
National Science Foundation
Michigan State University

The Essential Cosmic Perspective, Bennett, Donahue, Schneider, Voit, 7th Ed.

C. Wrede, HUGS, May 2019
Experimental Nuclear Astrophysics

Life of a massive star (more than about 8 solar masses)

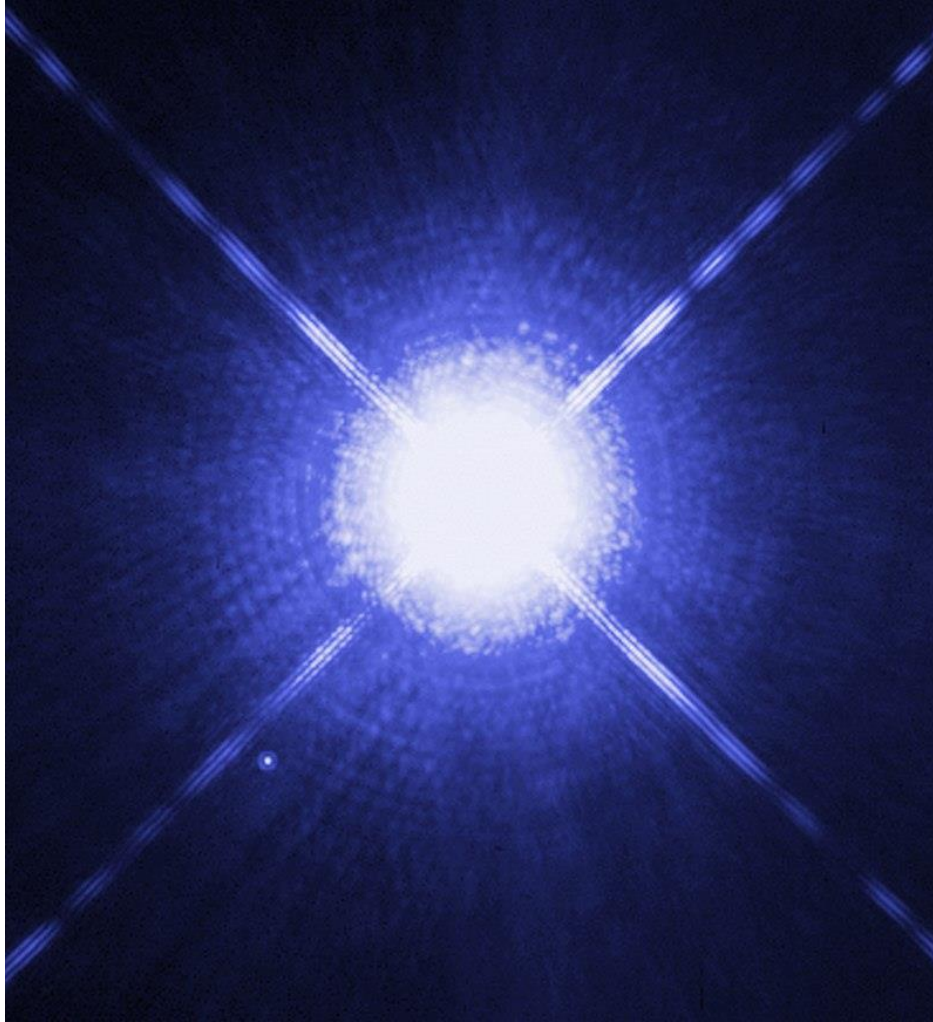
1. Protostar: cloud of cold gas collapses under gravity
2. Main Sequence: H fuses to He in core
3. Red Supergiant: H fuses to He in shell around inert He core
4. Helium Core Fusion: He fuses to C, O in core while H fuses to He in shell
5. Multiple Shell Fusion: many elements fuse in shells around Fe core
6. Supernova (type II) explosion after Fe core collapses
7. Neutron star or black hole left behind



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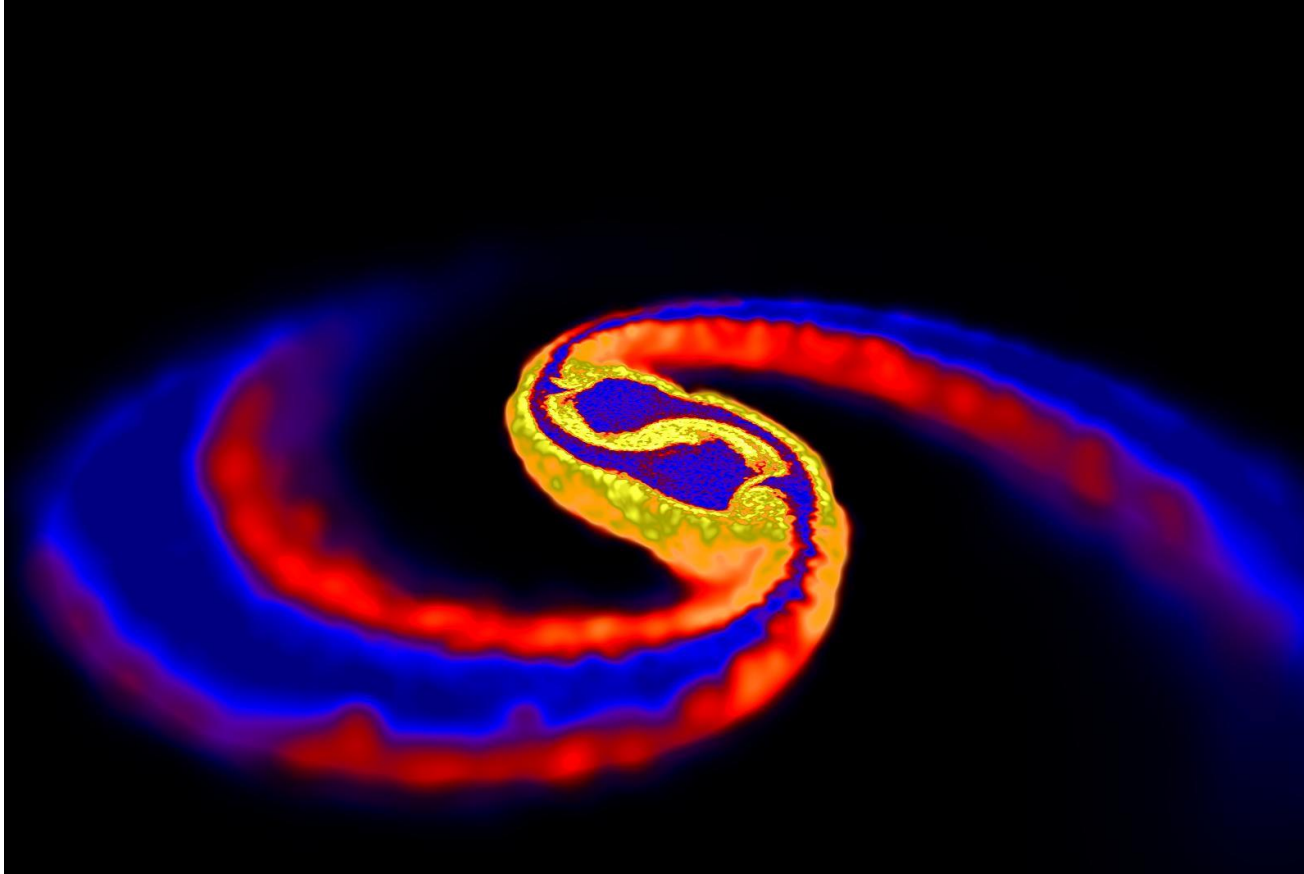
Binary star systems

Half of all stars are in binary systems. For example, the brightest star system in our sky Sirius.



NASA, ESA, H. Bond (STScI), and M. Barstow (University of Leicester)

Close binary star systems



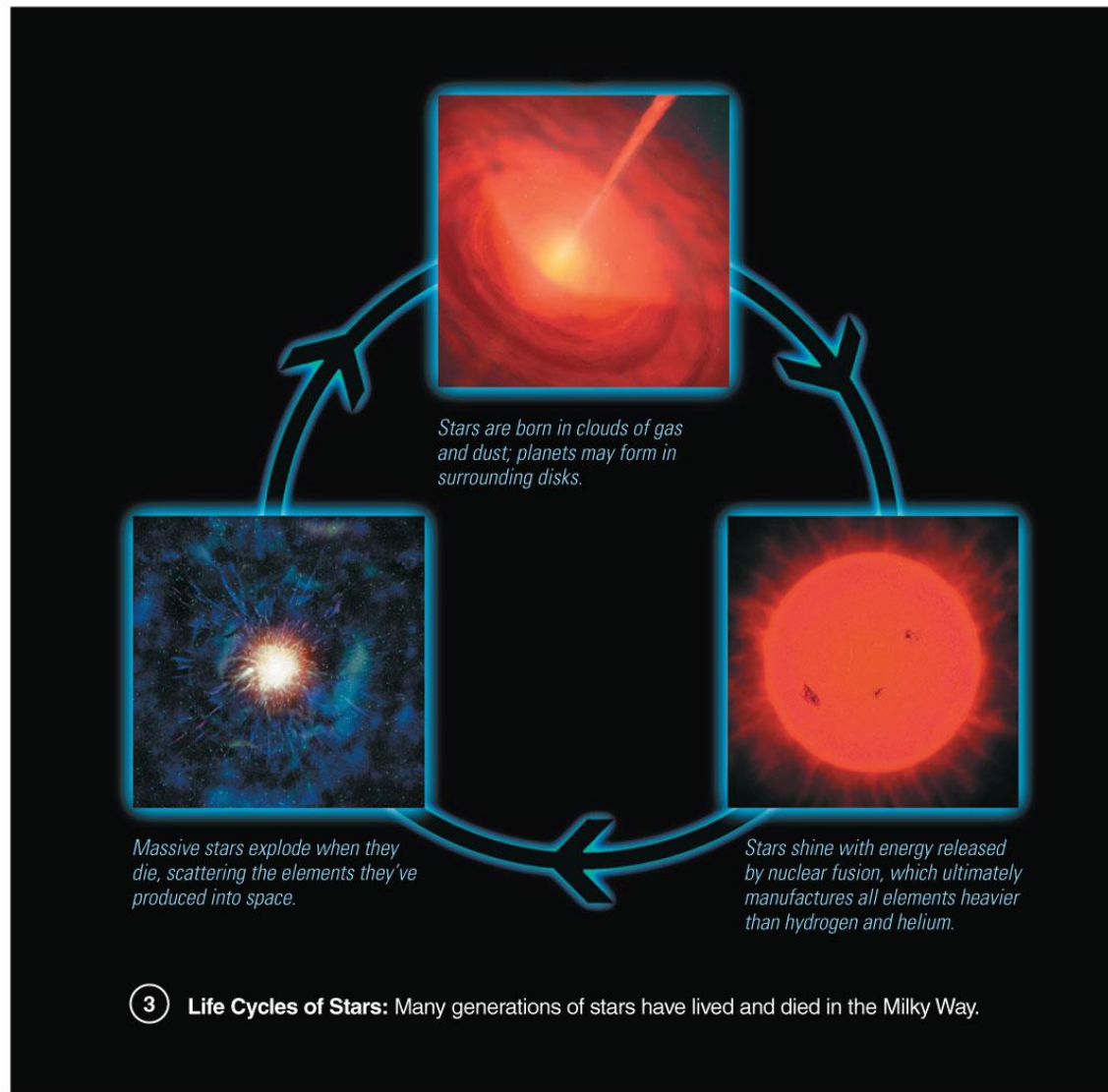
Interactions can lead to mass exchange and co-evolution, or rebirth of compact objects through neutron-star mergers, thermonuclear (Ia) supernovae, classical novae, X-ray bursts, ...

The Milky Way: a cosmic recycling plant



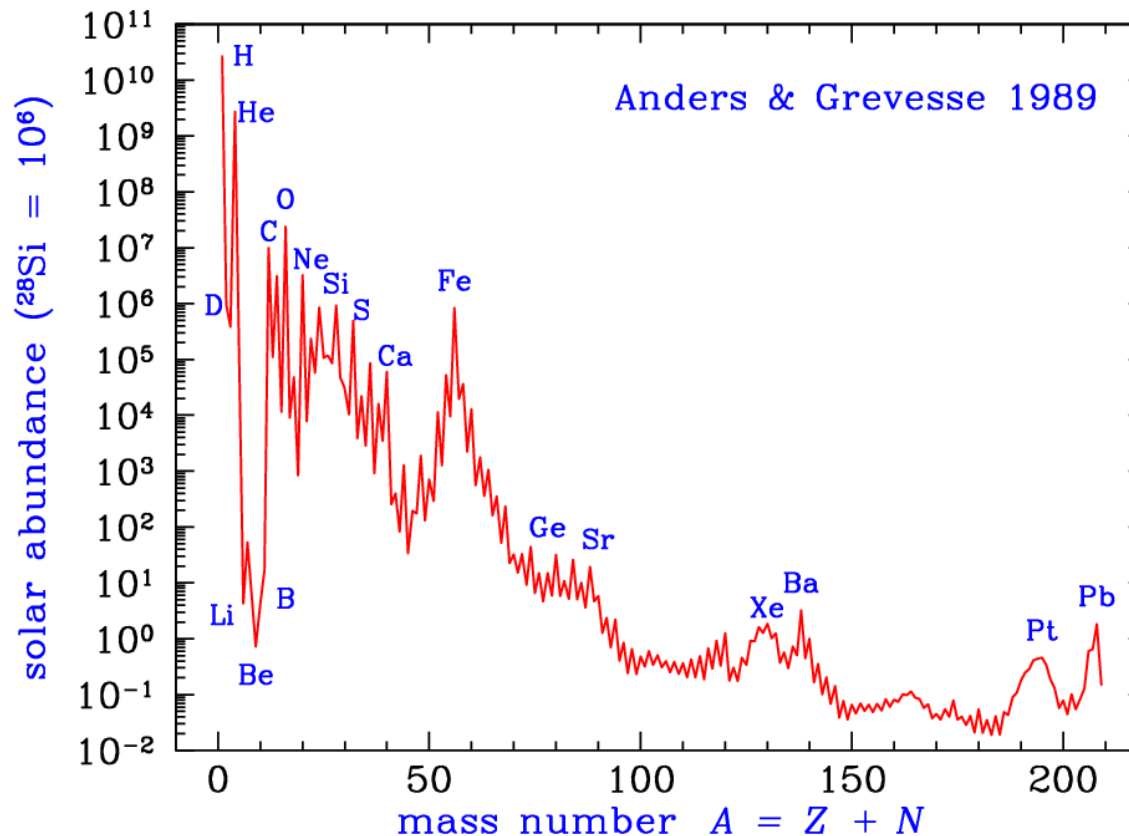
M31, the great galaxy in
Andromeda

The Milky Way: a cosmic recycling plant



Solar abundances

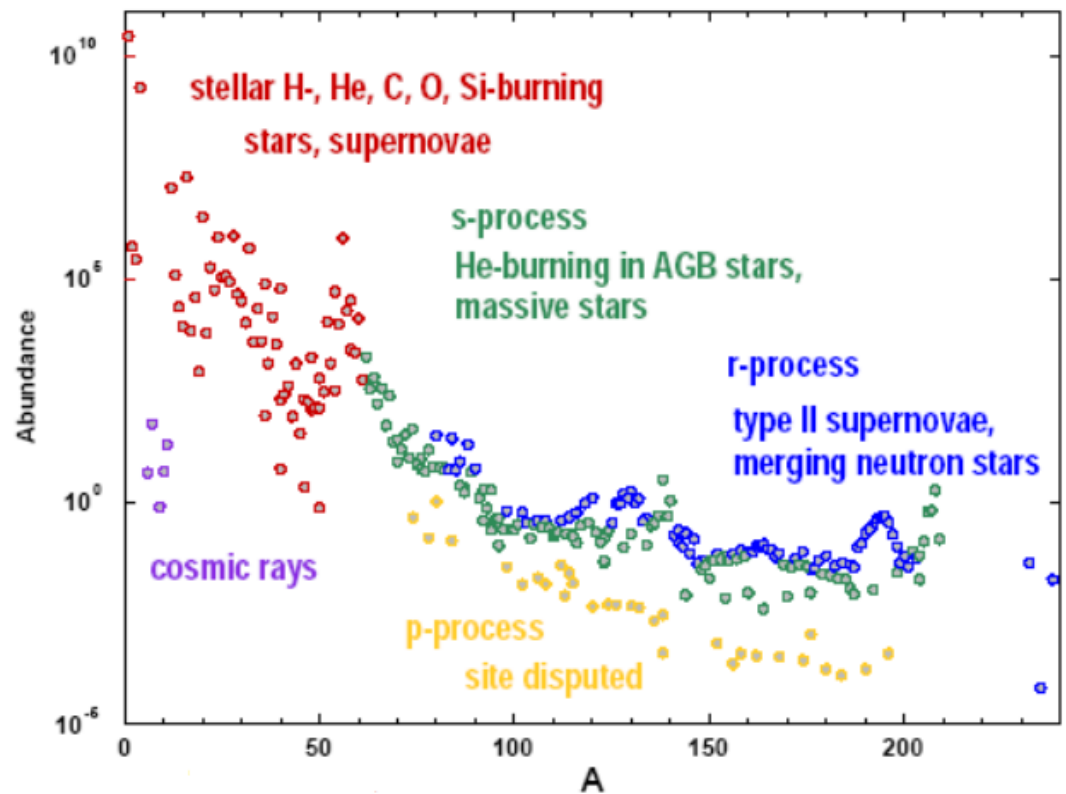
- Abundances determined by spectroscopy of the solar photosphere & mass spectrometry of primitive meteorites



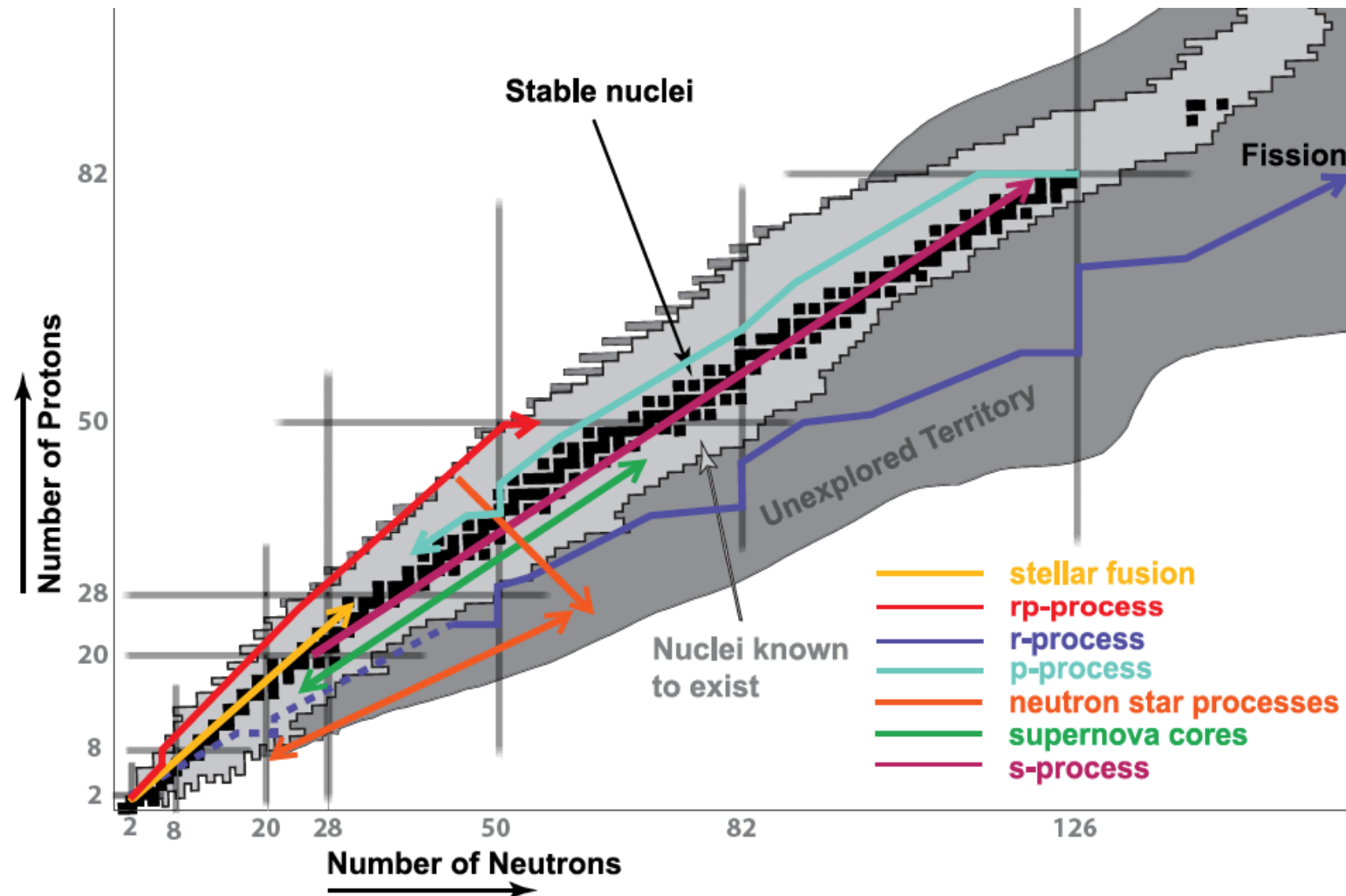
Synthesis of elements in stars

1. H burning \rightarrow conversion of H to He
2. He burning \rightarrow conversion of He to C, O ...
3. C, O and Ne burning \rightarrow production of A : 16 to 28
4. Si burning \rightarrow production of A : 28 to 60
5. s-, r- and p-processes \rightarrow production of $A > 60$
6. Li, Be, and B from cosmic rays

Solar abundances



Nuclear astrophysics processes



Need stellar nuclear reaction rates to understand nucleosynthesis processes

Nuclear Mass and Binding Energy

Nuclear mass:

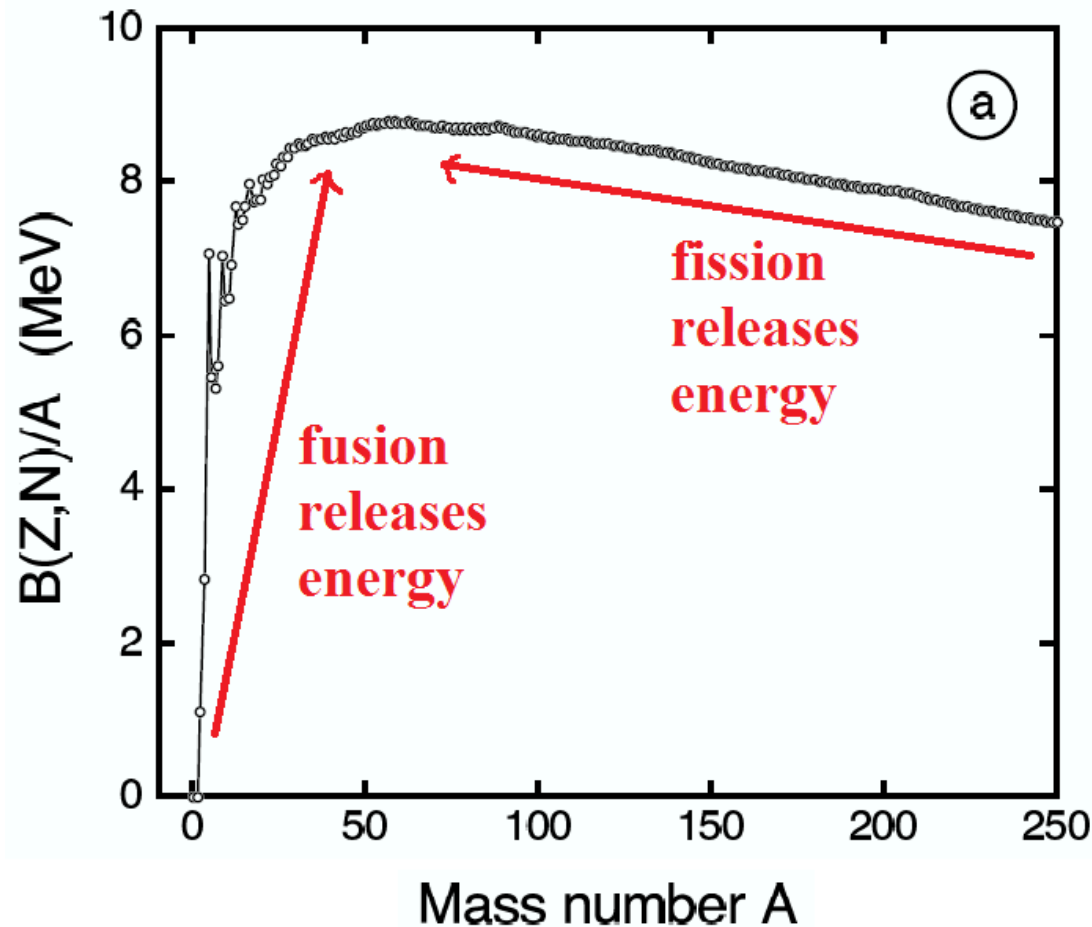
$$m_{\text{nuc}} = Zm_{\text{p}} + Nm_{\text{n}} - \Delta m$$

Nuclear binding energy:

$$\Delta m \cdot c^2 = \Delta E = B(Z, N)$$

$$B(Z, N) = (Zm_{\text{p}} + Nm_{\text{n}} - m_{\text{nuc}}) c^2$$

Binding energy per nucleon

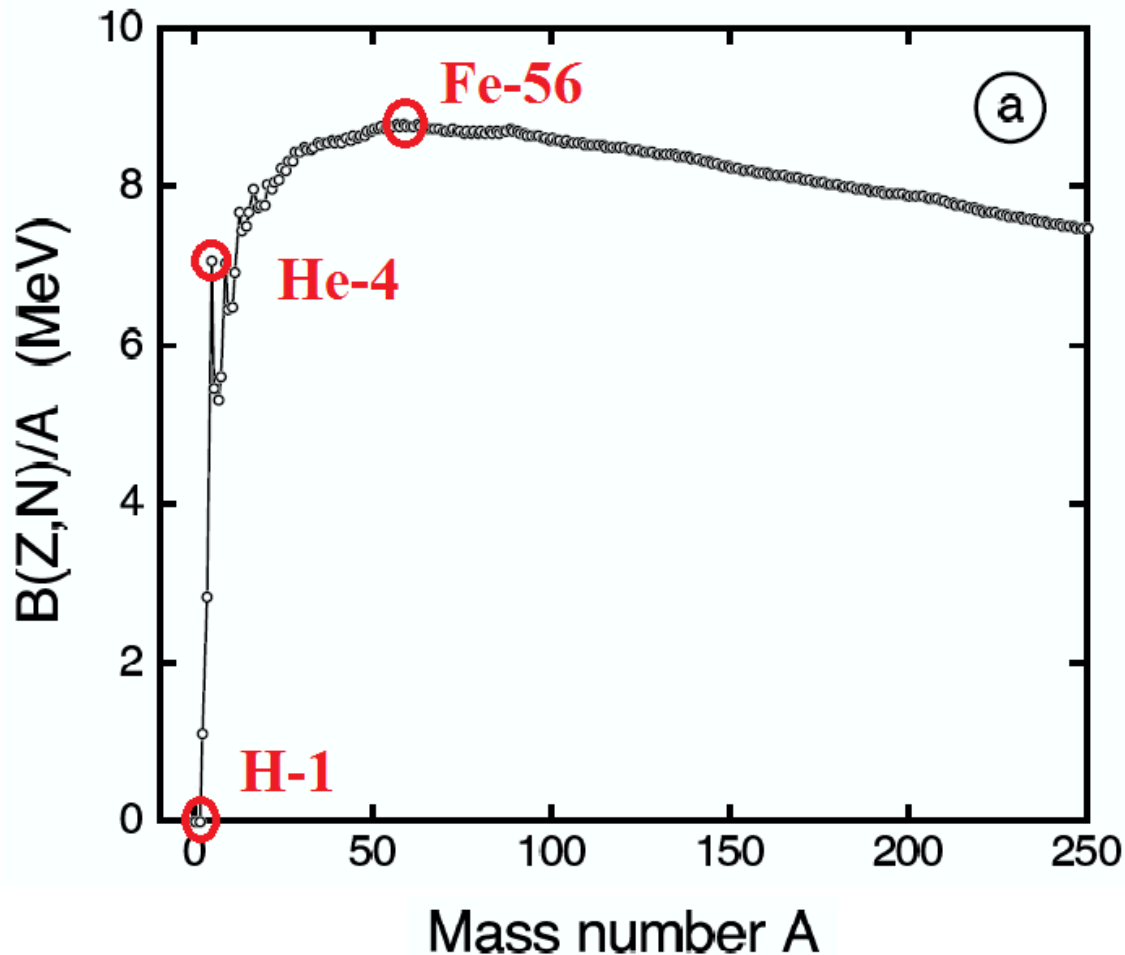


General trends:

Fusion of light nuclei releases energy; fission of heavy nuclei releases energy.

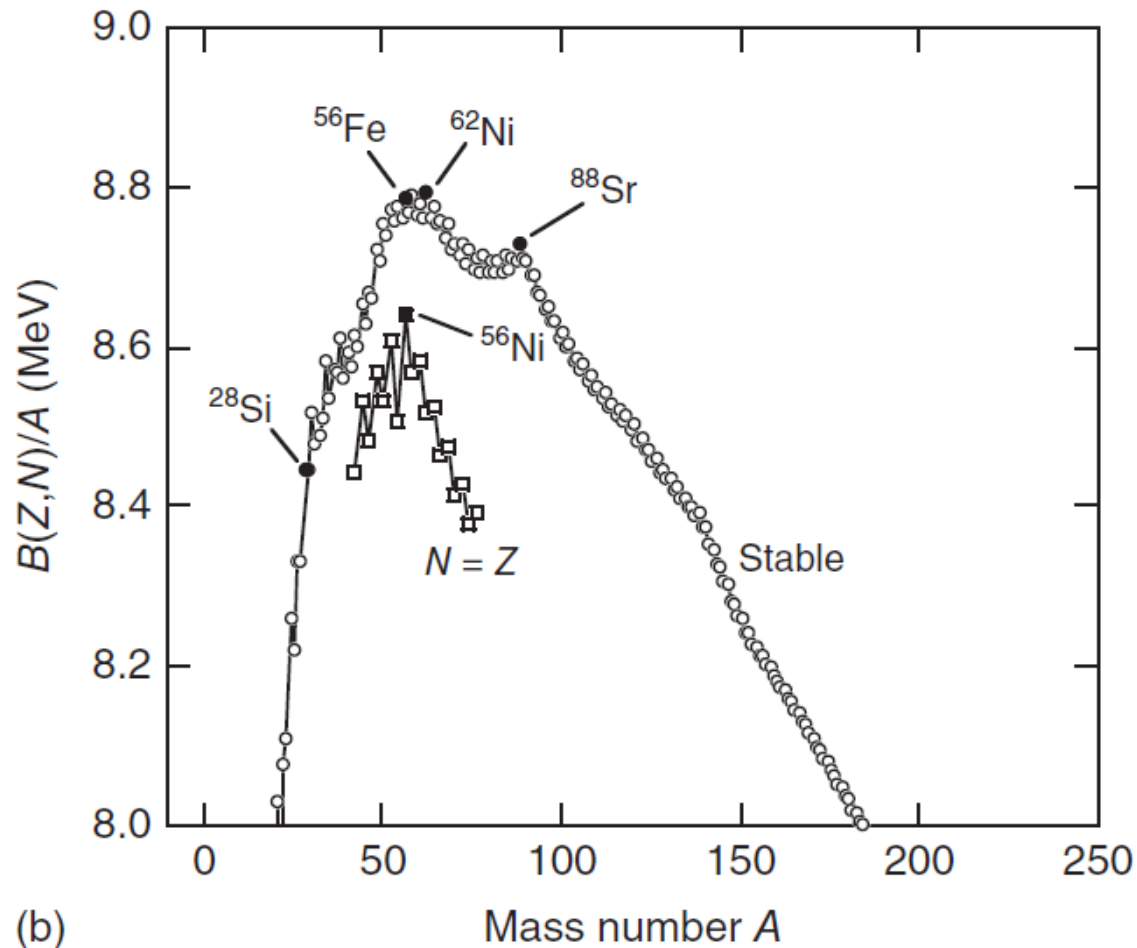
Fusion of heavy nuclei consumes energy; break-up of light nuclei consumes energy.

Binding energy per nucleon



Eg. fusing four ^1H into ^4He releases ~ 7 MeV per nucleon;
fusing many ^1H into ^{56}Fe releases ~ 9 MeV per nucleon.

Binding energy per nucleon



Peak in the range of $A = 50$ -65: iron peak nuclei! Nature favors production of the most tightly bound and stable nuclides. ^{62}Ni , ^{58}Fe , and ^{56}Fe are most bound.

Nuclear reactions: notation and terminology

A nuclear reaction can be represented symbolically as:

$$0 + 1 \rightarrow 2 + 3 \quad \text{or} \quad 0(1, 2)3$$

0,1 are colliding nuclei before reaction

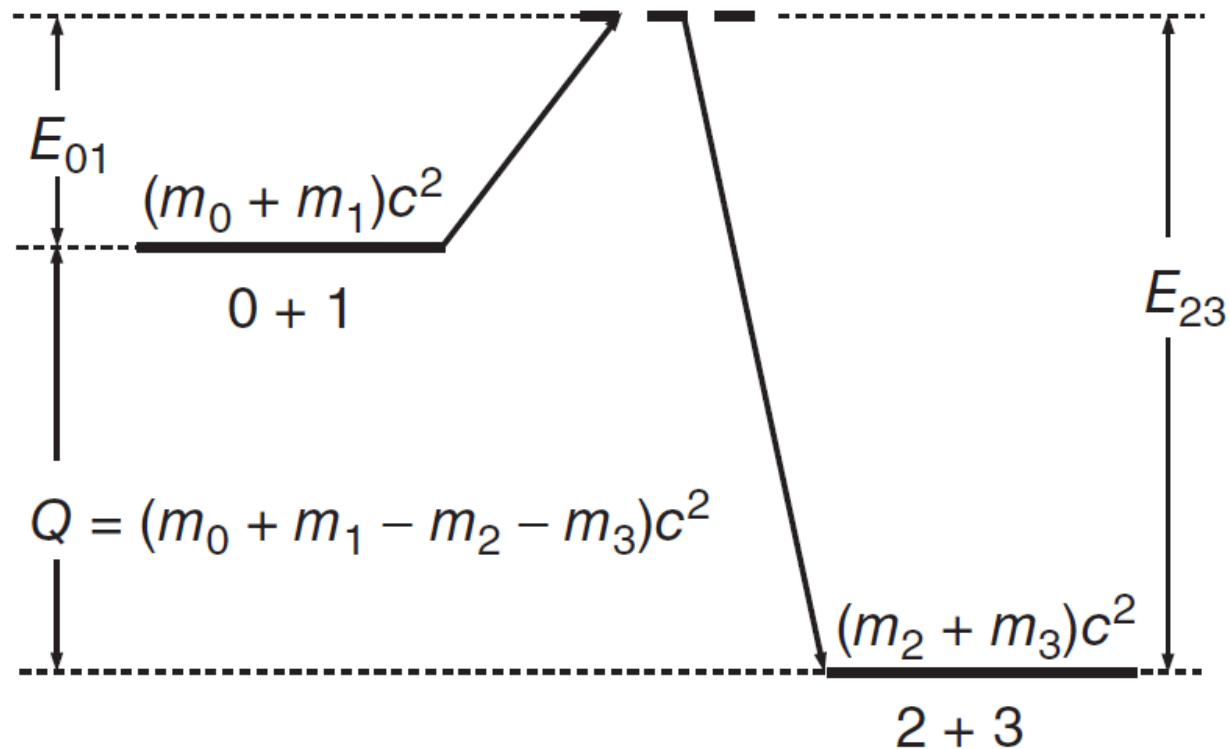
2,3 are the interaction products

If 0,1 identical to 2,3 then *elastic* or *inelastic scattering*;
otherwise *nuclear reaction*

If 2 is a photon: *radiative capture reaction*

If 1 is a photon: *photodisintegration reaction*

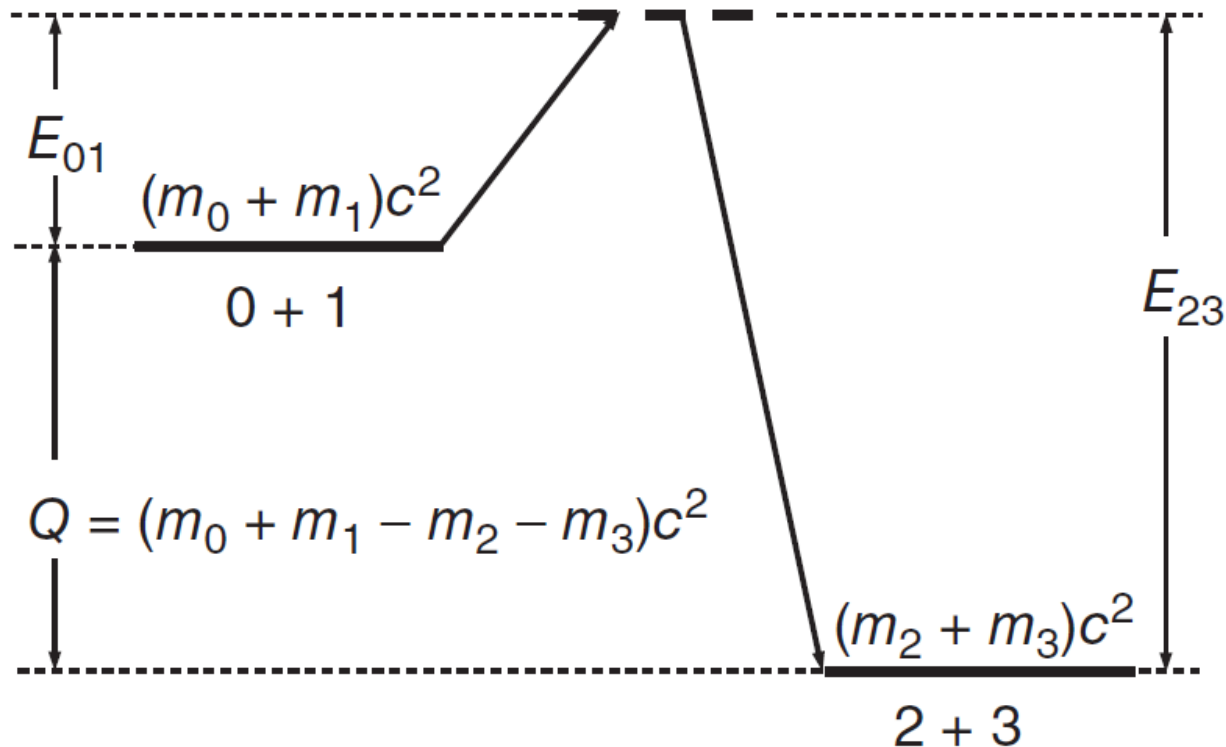
Energetics of nuclear reactions: 0(1,2)3



Conservation of energy:

$$m_0c^2 + m_1c^2 + E_0 + E_1 = m_2c^2 + m_3c^2 + E_2 + E_3$$

Energetics of nuclear reactions: 0(1,2)3



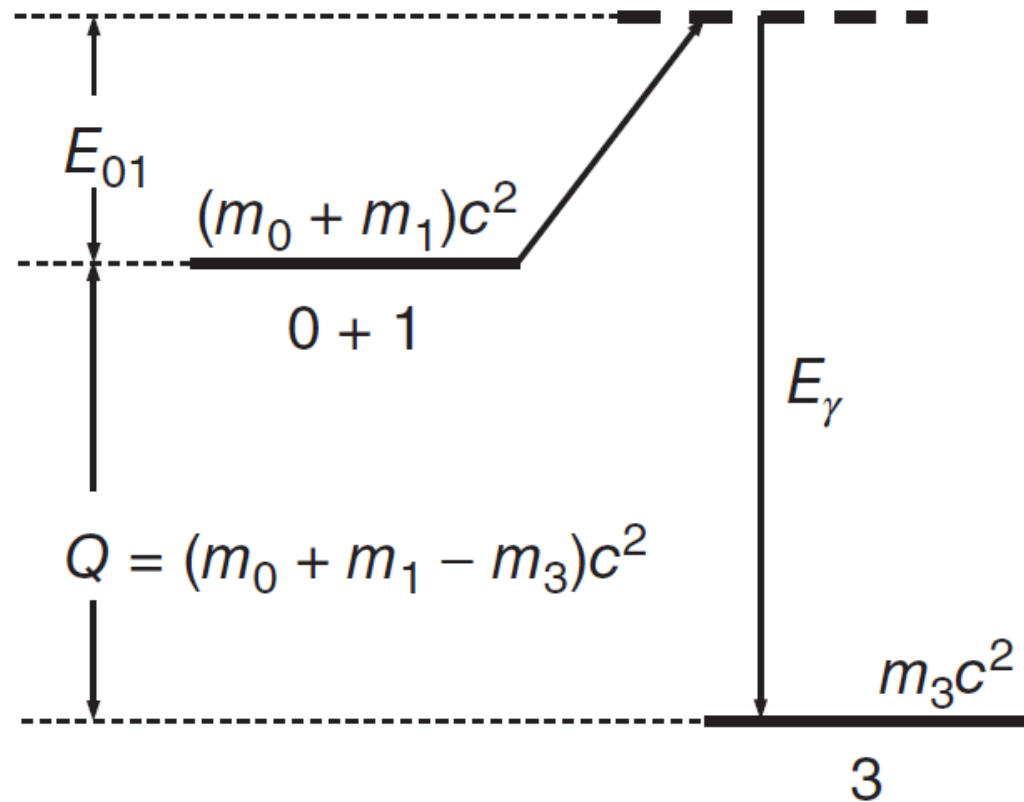
Definition of Q value:

$$Q_{01 \rightarrow 23} \equiv m_0 c^2 + m_1 c^2 - m_2 c^2 - m_3 c^2 = E_2 + E_3 - E_0 - E_1$$

If $Q > 0$: reaction releases energy (exothermic)

If $Q < 0$: reaction consumes energy (endothermic)

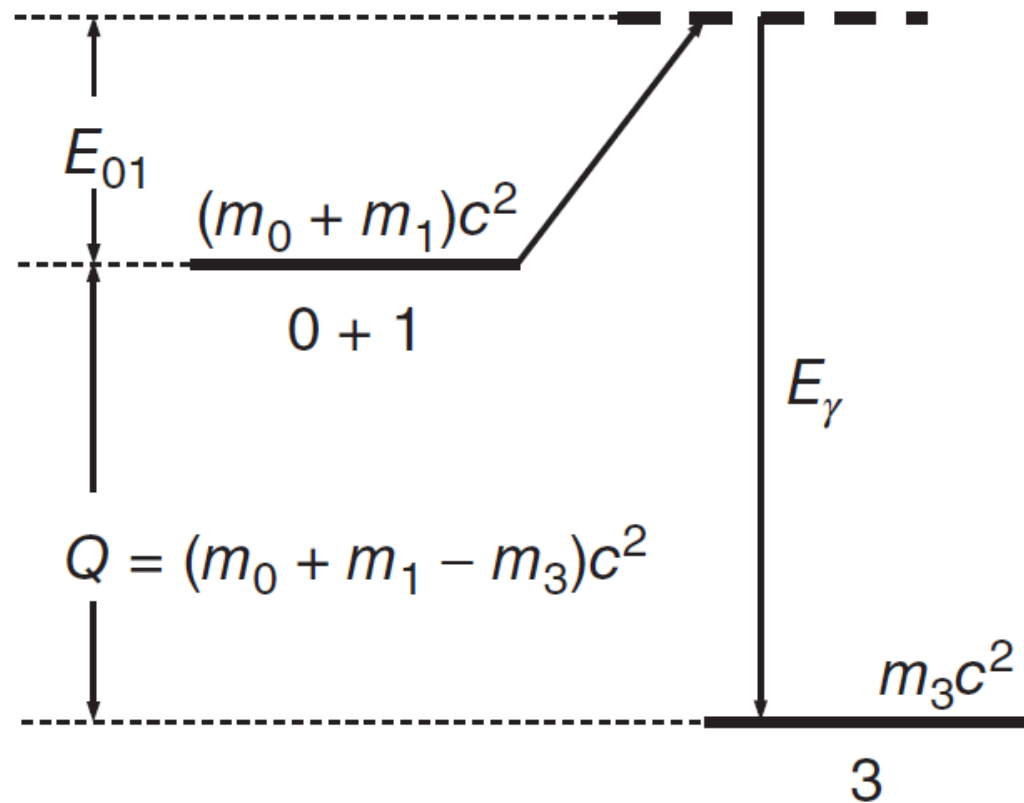
Radiative Capture Reactions: $0(1,\gamma)3$



$$m_0c^2 + m_1c^2 + E_0 + E_1 = m_3c^2 + E_3 + E_\gamma \quad \text{or}$$

$$Q_{01 \rightarrow \gamma 3} \equiv m_0c^2 + m_1c^2 - m_3c^2 = E_3 + E_\gamma - E_0 - E_1$$

Photodisintegration Reactions: $3(\gamma,1)0$



If one adds energy $> Q$ to nucleus 3, then it becomes energetically possible for nucleus 3 to separate into fragments 0,1. Define *particle separation energy*:

$$S_{3 \rightarrow 01} = Q_{01 \rightarrow \gamma 3}$$

Stellar reaction rates: *thermonuclear*

stellar reaction rate $\langle \sigma v \rangle = \int \sigma(v) \phi(v) v dv$

- need:
- a) velocity distribution $\phi(v)$
 - b) cross section $\sigma(v)$

a) velocity distribution

interacting nuclei in plasma are in **thermal equilibrium** at temperature T

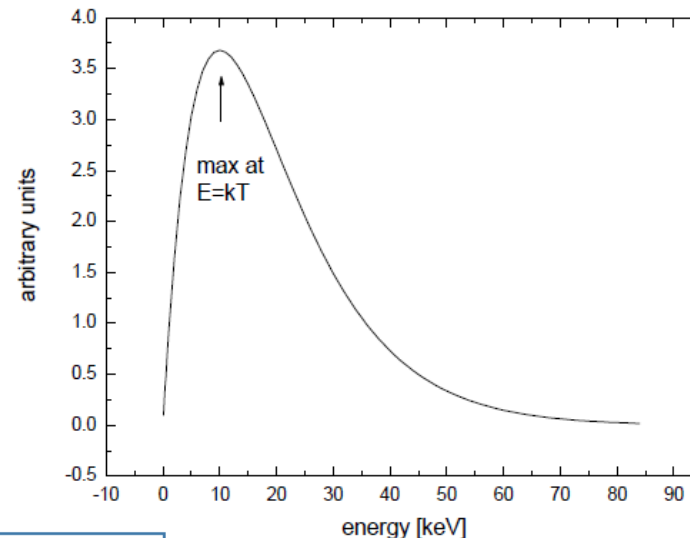
also assume **non-degenerate** and **non-relativistic** plasma

⇒ **Maxwell-Boltzmann velocity distribution**

$$\phi(v) = 4\pi \left(\frac{\mu}{2\pi kT} \right)^{3/2} v^2 \exp\left(-\frac{\mu v^2}{2kT} \right)$$

with $\mu = \frac{m_p m_T}{m_p + m_T}$ reduced mass

v = relative velocity



$$kT \sim 8.6 \times 10^{-8} T[\text{K}] \text{ keV}$$

example: Sun $T \sim 15 \times 10^6 \text{ K}$ ⇒ $kT \sim 1 \text{ keV}$

Stellar reaction rates: thermonuclear

b) cross section

Probability for reaction between nuclei (units of area)

no nuclear theory available to determine reaction cross section a priori

cross section depends sensitively on:

- the properties of the nuclei involved
- the reaction mechanism

and can vary by orders of magnitude, depending on the interaction

examples:

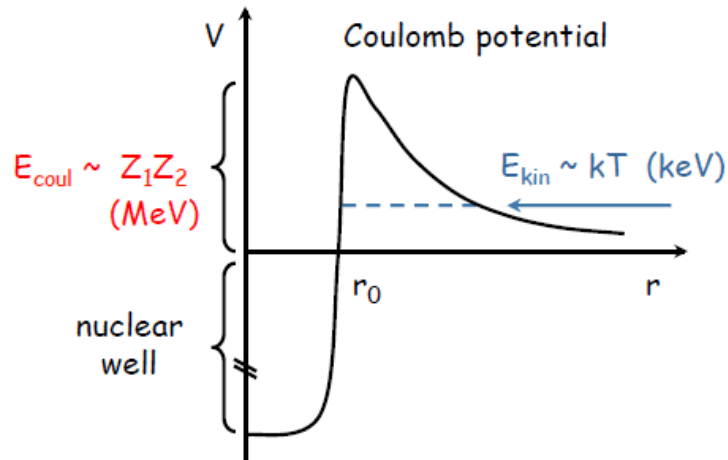
Reaction	Force	σ (barn)	E_{proj} (MeV)
$^{15}\text{N}(p, \alpha)^{12}\text{C}$	strong	0.5	2.0
$^3\text{He}(\alpha, \gamma)^7\text{Be}$	electromagnetic	10^{-6}	2.0
$p(p, e^+ \nu)d$	weak	10^{-20}	2.0

$$1 \text{ barn} = 10^{-24} \text{ cm}^2 = 100 \text{ fm}^2$$

Stellar reaction rates: *thermonuclear* (charged particles)

charged particles → Coulomb barrier

energy available: from thermal motion

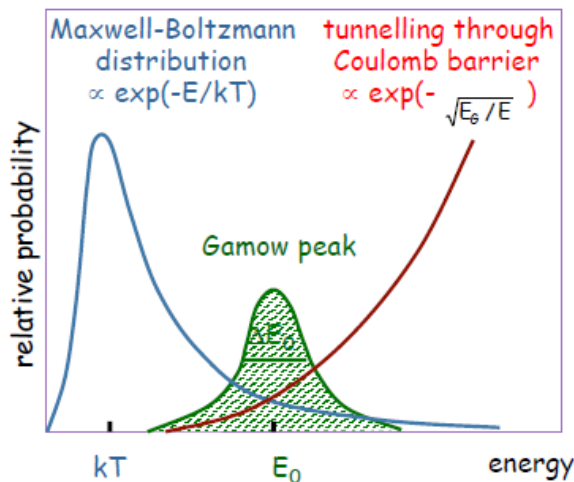


during static burning: $kT \ll E_{\text{coul}}$

$T \sim 15 \times 10^6 \text{ K}$ (e.g. our Sun) $\Rightarrow kT \sim 1 \text{ keV}$

reactions occur through **TUNNEL EFFECT**

→ tunneling probability $P \propto \exp(-2\pi\eta)$



Gamow peak: energy of astrophysical interest
where measurements should be carried out

$$kT \ll E_0 \ll E_{\text{coul}}$$

$$10^{-18} \text{ barn} < \sigma < 10^{-9} \text{ barn}$$



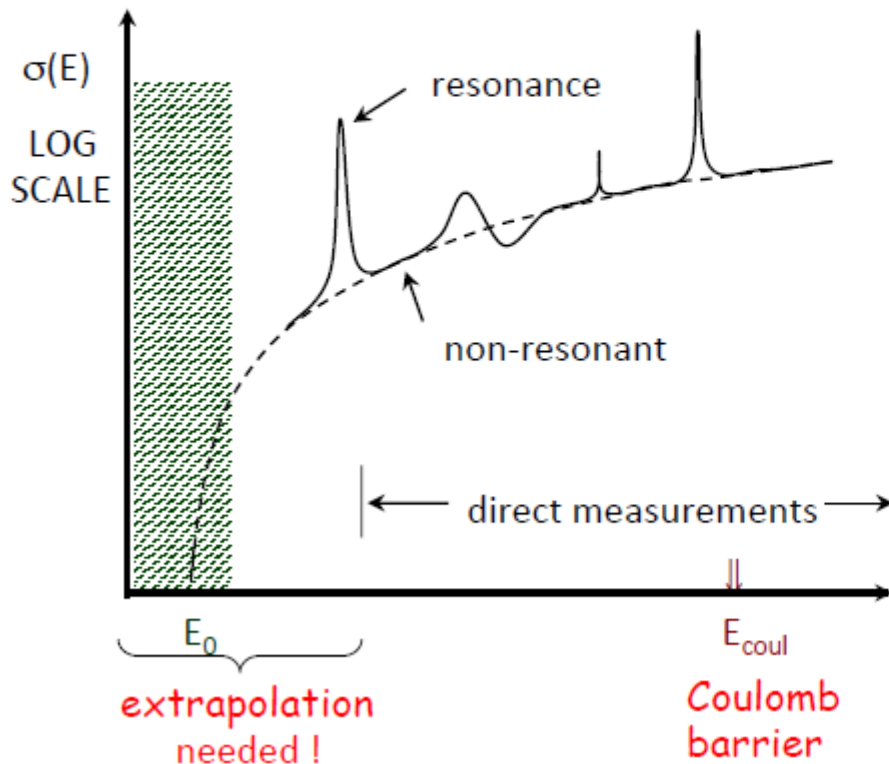
major experimental challenges

Experimental approach (charged particles)

Measure $\sigma(E)$ as low as possible in energy and extrapolate if necessary

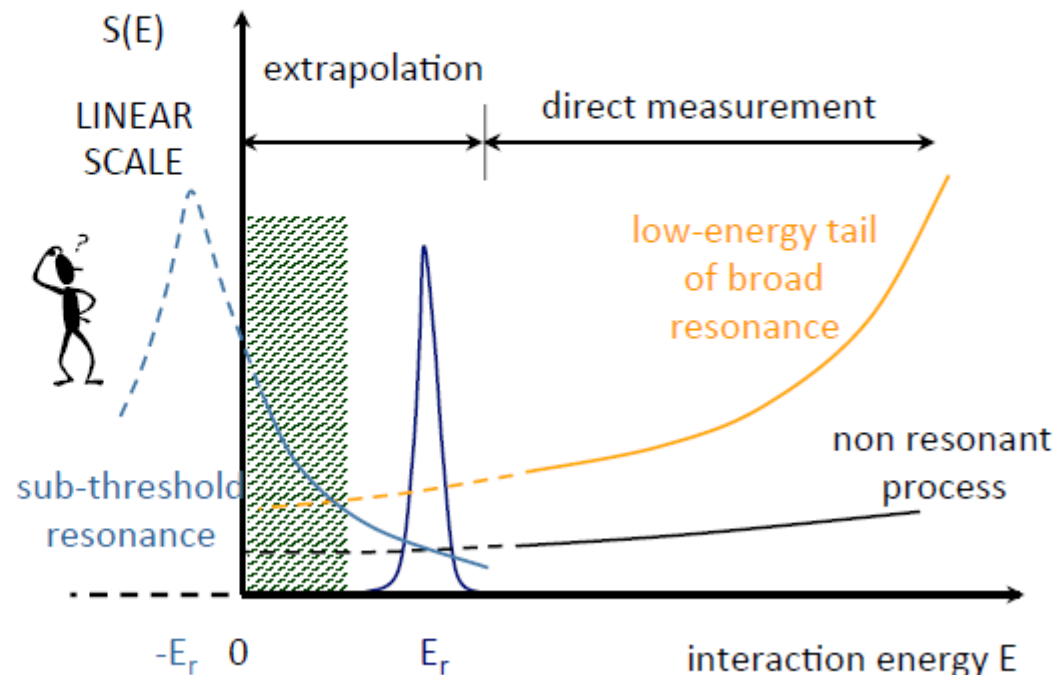
CROSS SECTION

$$\sigma(E) = \frac{1}{E} \exp(-2\pi\eta) S(E)$$



S-FACTOR

$$S(E) = E \sigma(E) \exp(2\pi\eta)$$



DANGER OF EXTRAPOLATION !

Neutron capture reaction rates

- No Coulomb barrier, so cross sections are larger at low energies

$$\langle \sigma v \rangle = \int \sigma(v) \phi(v) v dv = \int \sigma(E) \exp(-E/kT) E dE$$

s-wave neutron capture

energy range of interest $E \sim kT$

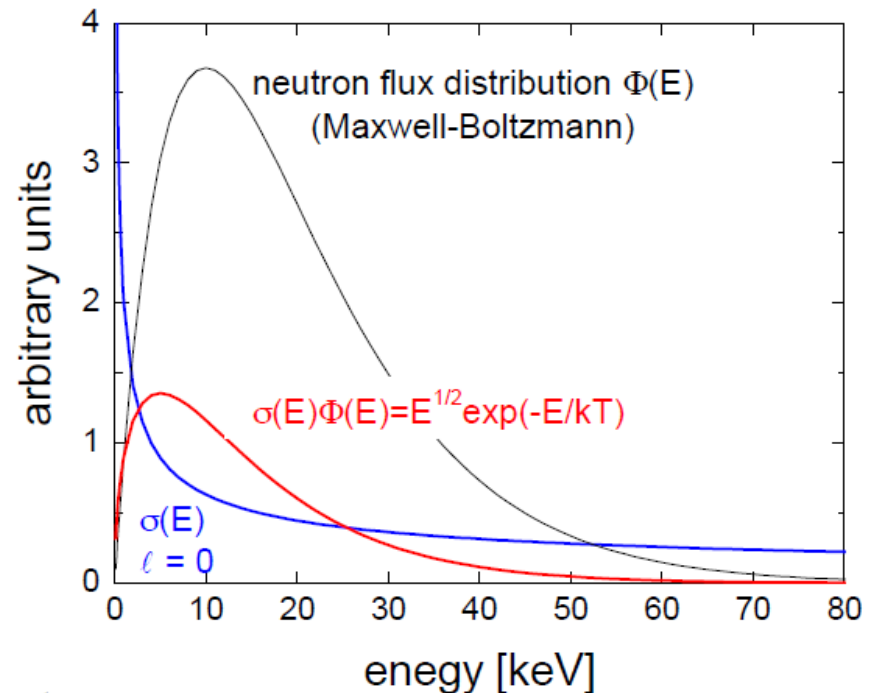
$$\sigma \propto \frac{1}{v} \Rightarrow \sigma v = \text{const} = \langle \sigma v \rangle$$

stellar reaction rate

$$\langle \sigma v \rangle = v_T \sigma_{\text{th}}$$

σ_{th} = measured cross section for thermal neutrons

$$v_T = \sqrt{\frac{2kT}{\mu}} \quad \text{most probable velocity, corresponding to } E_{\text{cm}} = kT$$



- Need low-energy neutron beams to bombard stable targets and measure yield

Summary

- Historical development of nuclear astrophysics
- Stellar evolution and galactic chemical evolution
- Thermonuclear reaction rates
- Next: direct measurements of charged-particle reactions

Thank you for your attention!